The algebraic and geometric classification of nilpotent left symmetric algebras ¹

Jobir Adashev^a, Ivan Kaygorodov^b, Abror Khudoyberdiyev^{a,c} & Aloberdi Sattarov^a

- ^a Institute of Mathematics Academy of Sciences of Uzbekistan, Tashkent, Uzbekistan.
- ^b CMCC, Universidade Federal do ABC, Santo André, Brazil.
- ^c National University of Uzbekistan, Tashkent, Uzbekistan.

E-mail addresses:

Jobir Adashev (adashevjq@mail.ru) Ivan Kaygorodov (kaygorodov.ivan@gmail.com) Abror Khudoyberdiyev (khabror@mail.ru) Aloberdi Sattarov (saloberdi@mail.ru)

Abstract: This paper is devoted to the complete algebraic and geometric classification of complex 4-dimensional nilpotent left symmetric algebras. It was proved that the variety is 15-dimensional and it has three irreducible components.

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Introduction

The algebraic classification (up to isomorphism) of algebras of dimension n from a certain variety defined by a certain family of polinomial identities is a classic problem in the theory of non-associative algebras. There are many results related to the algebraic classification of small-dimensional algebras in the varieties of Jordan, Lie, Leibniz, Zinbiel and many other algebras [1,11,14,17-19,23,27,29-31,35,40,43]. Another interesting direction in the classification of algebras is the geometric classification. There are many results related to the geometric classification of Jordan, Lie, Leibniz, Zinbiel and many other algebras [7,10,12,24-26,30-33,36,38-42,46,49]. An algebraic classification of complex 3-dimensional left symmetric algebras is given in [4]. In the present paper, we give the algebraic and geometric classification of 4-dimensional nilpotent left symmetric algebras.

Left-symmetric algebras (or under other names like KoszulVinberg algebras, quasi-associative algebras, pre-Lie algebras, and so on) are a class of nonassociative algebras coming from the study of several topics in geometry and algebra, such as rooted tree algebras [13], convex homogeneous cones [48], affine manifolds and affine structures on Lie groups [44], deformation of associative algebras [22], and so on. They are Lie-admissible algebras (in the sense that the commutators define Lie algebra structures) whose left multiplication operators form a Lie algebra. It contains associative algebras, Novikov algebras and assosymmetric algebras as subvarieties. The variety of left symmetric algebras is defined by the following identity:

$$(xy)z - x(yz) = (yx)z - y(xz).$$

Furthermore, left-symmetric algebras are a kind of natural algebraic systems appearing in many fields in mathematics and mathematical physics. Perhaps this is one of the most attractive and interesting places. As it was pointed out in a paper of Chapoton and Livernet [15], the left-symmetric algebra deserves more attention than it has been given. For example, left-symmetric algebras appear as an underlying structure of those Lie algebras that possess a phase space, thus they form a natural category from the point of view of classical and quantum mechanics [45]; they are the underlying algebraic structures of vertex algebras [5]; there is a correspondence between left-symmetric algebras and complex product structures on Lie algebras [3], which plays an important role in the theory of hypercomplex and hypersymplectic manifolds;

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left-symmetric algebras have close relations with certain integrable systems [8], classical and quantum YangBaxter equation [20]; Poisson brackets and infinite-dimensional Lie algebras [6, 21], operads [15], quantum field theory [16], and so on (see [9] and the references therein).

Our method for classifying nilpotent left symmetric algebras is based on the calculation of central extensions of nilpotent algebras of smaller dimensions from the same variety. The algebraic study of central extensions of Lie and non-Lie algebras has been an important topic for years [2, 28, 47, 50]. First, Skjelbred and Sund used central extensions of Lie algebras to obtain a classification of nilpotent Lie algebras [47]. After that, using the method described by Skjelbred and Sund, all non-Lie central extensions of all 4-dimensional Malcev algebras were described [28], and also all non-associative central extensions of 3-dimensional Jordan algebras. Note that the Skjelbred-Sund method of central extensions is an important tool in the classification of nilpotent algebras, which was used to describe all 4-dimensional nilpotent associative algebras [19], all 4-dimensional nilpotent bicommutative algebras [37], all 5-dimensional nilpotent Jordan algebras [27], all 5-dimensional nilpotent restricted Lie algebras [17], all 6-dimensional nilpotent Lie algebras [18], all 6-dimensional nilpotent Malcev algebras [29] and some others.

1. THE ALGEBRAIC CLASSIFICATION OF NILPOTENT LEFT SYMMETRIC ALGEBRAS

1.1. **Method of classification of nilpotent algebras.** Throughout this paper, we use the notations and methods well written in [28], which we have adapted for the left symmetric case with some modifications. Further in this section we give some important definitions.

Let (\mathbf{A},\cdot) be a left symmetric algebra over $\mathbb C$ and $\mathbb V$ a vector space over $\mathbb C$. The $\mathbb C$ -linear space $Z^2(\mathbf{A},\mathbb V)$ is defined as the set of all bilinear maps $\theta\colon \mathbf{A}\times \mathbf{A}\longrightarrow \mathbb V$ such that

$$\theta(xy, z) - \theta(x, yz) = \theta(yx, z) - \theta(y, xz).$$

These elements will be called *cocycles*. For a linear map f from \mathbf{A} to \mathbb{V} , if we define $\delta f \colon \mathbf{A} \times \mathbf{A} \longrightarrow \mathbb{V}$ by $\delta f(x,y) = f(xy)$, then $\delta f \in \mathrm{Z}^2(\mathbf{A},\mathbb{V})$. We define $\mathrm{B}^2(\mathbf{A},\mathbb{V}) = \{\theta = \delta f : f \in \mathrm{Hom}(\mathbf{A},\mathbb{V})\}$. We define the *second cohomology space* $\mathrm{H}^2(\mathbf{A},\mathbb{V})$ as the quotient space $\mathrm{Z}^2(\mathbf{A},\mathbb{V})/\mathrm{B}^2(\mathbf{A},\mathbb{V})$.

Let $\operatorname{Aut}(\mathbf{A})$ be the automorphism group of \mathbf{A} and let $\phi \in \operatorname{Aut}(\mathbf{A})$. For $\theta \in \operatorname{Z}^2(\mathbf{A}, \mathbb{V})$ define the action of the group $\operatorname{Aut}(\mathbf{A})$ on $\operatorname{Z}^2(\mathbf{A}, \mathbb{V})$ by $\phi\theta(x,y) = \theta(\phi(x), \phi(y))$. It is easy to verify that $\operatorname{B}^2(\mathbf{A}, \mathbb{V})$ is invariant under the action of $\operatorname{Aut}(\mathbf{A})$. So, we have an induced action of $\operatorname{Aut}(\mathbf{A})$ on $\operatorname{H}^2(\mathbf{A}, \mathbb{V})$.

Let **A** be a left symmetric algebra of dimension m over \mathbb{C} and \mathbb{V} be a \mathbb{C} -vector space of dimension k. For the bilinear map θ , define on the linear space $\mathbf{A}_{\theta} = \mathbf{A} \oplus \mathbb{V}$ the bilinear product " $[-,-]_{\mathbf{A}_{\theta}}$ " by $[x+x',y+y']_{\mathbf{A}_{\theta}} = xy + \theta(x,y)$ for all $x,y \in \mathbf{A}, x',y' \in \mathbb{V}$. The algebra \mathbf{A}_{θ} is called a k-dimensional central extension of \mathbf{A} by \mathbb{V} . One can easily check that \mathbf{A}_{θ} is a left symmetric algebra if and only if $\theta \in \mathbb{Z}^2(\mathbf{A}, \mathbb{V})$.

Call the set $\operatorname{Ann}(\theta) = \{x \in \mathbf{A} : \theta(x, \mathbf{A}) + \theta(\mathbf{A}, x) = 0\}$ the *annihilator* of θ . We recall that the *annihilator* of an algebra \mathbf{A} is defined as the ideal $\operatorname{Ann}(\mathbf{A}) = \{x \in \mathbf{A} : x\mathbf{A} + \mathbf{A}x = 0\}$. Observe that $\operatorname{Ann}(\mathbf{A}_{\theta}) = (\operatorname{Ann}(\theta) \cap \operatorname{Ann}(\mathbf{A})) \oplus \mathbb{V}$.

The following result shows that every algebra with a non-zero annihilator is a central extension of a smaller-dimensional algebra.

Lemma 1. Let A be an n-dimensional left symmetric algebra such that $\dim(\operatorname{Ann}(A)) = m \neq 0$. Then there exists, up to isomorphism, a unique (n-m)-dimensional left symmetric algebra A' and a bilinear map $\theta \in Z^2(A', \mathbb{V})$ with $\operatorname{Ann}(A') \cap \operatorname{Ann}(\theta) = 0$, where \mathbb{V} is a vector space of dimension m, such that $A \cong A'_{\theta}$ and $A / \operatorname{Ann}(A) \cong A'$.

Proof. Let A' be a linear complement of Ann(A) in A. Define a linear map $P: A \longrightarrow A'$ by P(x+v) = x for $x \in A'$ and $v \in Ann(A)$, and define a multiplication on A' by $[x,y]_{A'} = P(xy)$ for $x,y \in A'$. For $x,y \in A$, we have

$$P(xy) = P((x - P(x) + P(x))(y - P(y) + P(y))) = P(P(x)P(y)) = [P(x), P(y)]_{\mathbf{A}'}.$$

Since P is a homomorphism $P(\mathbf{A}) = \mathbf{A}'$ is a left symmetric algebra and $\mathbf{A}/\operatorname{Ann}(\mathbf{A}) \cong \mathbf{A}'$, which gives us the uniqueness. Now, define the map $\theta \colon \mathbf{A}' \times \mathbf{A}' \longrightarrow \operatorname{Ann}(\mathbf{A})$ by $\theta(x,y) = xy - [x,y]_{\mathbf{A}'}$. Thus, \mathbf{A}'_{θ} is \mathbf{A} and therefore $\theta \in Z^2(\mathbf{A}', \mathbb{V})$ and $\operatorname{Ann}(\mathbf{A}') \cap \operatorname{Ann}(\theta) = 0$.

Definition 2. Let A be an algebra and I be a subspace of Ann(A). If $A = A_0 \oplus I$ then I is called an annihilator component of A. A central extension of an algebra A without annihilator component is called a non-split central extension.

Our task is to find all central extensions of an algebra A by a space V. In order to solve the isomorphism problem we need to study the action of Aut(A) on $H^2(A, V)$. To do that, let us fix a basis e_1, \ldots, e_s of

$$\mathbb{V}$$
, and $\theta \in \mathbb{Z}^{2}(\mathbf{A}, \mathbb{V})$. Then θ can be uniquely written as $\theta(x, y) = \sum_{i=1}^{s} \theta_{i}(x, y) e_{i}$, where $\theta_{i} \in \mathbb{Z}^{2}(\mathbf{A}, \mathbb{C})$.

Moreover, $\operatorname{Ann}(\theta) = \operatorname{Ann}(\theta_1) \cap \operatorname{Ann}(\theta_2) \cap \ldots \cap \operatorname{Ann}(\theta_s)$. Furthermore, $\theta \in \operatorname{B}^2(\mathbf{A}, \mathbb{V})$ if and only if all $\theta_i \in \operatorname{B}^2(\mathbf{A}, \mathbb{C})$. It is not difficult to prove (see [28, Lemma 13]) that given a left symmetric algebra \mathbf{A}_{θ} ,

if we write as above $\theta(x,y) = \sum_{i=1}^{s} \theta_i(x,y) e_i \in \mathrm{Z}^2(\mathbf{A},\mathbb{V})$ and $\mathrm{Ann}(\theta) \cap \mathrm{Ann}(\mathbf{A}) = 0$, then \mathbf{A}_{θ} has an annihilator component if and only if $[\theta_1], [\theta_2], \dots, [\theta_s]$ are linearly dependent in $\mathrm{H}^2(\mathbf{A},\mathbb{C})$.

Let \mathbb{V} be a finite-dimensional vector space over \mathbb{C} . The *Grassmannian* $G_k(\mathbb{V})$ is the set of all k-dimensional linear subspaces of \mathbb{V} . Let $G_s(H^2(\mathbf{A},\mathbb{C}))$ be the Grassmannian of subspaces of dimension s in $H^2(\mathbf{A},\mathbb{C})$. There is a natural action of $\operatorname{Aut}(\mathbf{A})$ on $G_s(H^2(\mathbf{A},\mathbb{C}))$. Let $\phi \in \operatorname{Aut}(\mathbf{A})$. For $W = \langle [\theta_1], [\theta_2], \ldots, [\theta_s] \rangle \in G_s(H^2(\mathbf{A},\mathbb{C}))$ define $\phi W = \langle [\phi \theta_1], [\phi \theta_2], \ldots, [\phi \theta_s] \rangle$. We denote the orbit of $W \in G_s(H^2(\mathbf{A},\mathbb{C}))$ under the action of $\operatorname{Aut}(\mathbf{A})$ by $\operatorname{Orb}(W)$. Given

$$W_1 = \langle [\theta_1], [\theta_2], \dots, [\theta_s] \rangle, W_2 = \langle [\vartheta_1], [\vartheta_2], \dots, [\vartheta_s] \rangle \in G_s \left(H^2 \left(\mathbf{A}, \mathbb{C} \right) \right),$$

we easily have that if $W_1 = W_2$, then $\bigcap_{i=1}^s \operatorname{Ann}(\theta_i) \cap \operatorname{Ann}(\mathbf{A}) = \bigcap_{i=1}^s \operatorname{Ann}(\vartheta_i) \cap \operatorname{Ann}(\mathbf{A})$, and therefore we can introduce the set

$$\mathbf{T}_{s}(\mathbf{A}) = \left\{ W = \langle [\theta_{1}], [\theta_{2}], \dots, [\theta_{s}] \rangle \in G_{s} \left(\mathbf{H}^{2} \left(\mathbf{A}, \mathbb{C} \right) \right) : \bigcap_{i=1}^{s} \operatorname{Ann}(\theta_{i}) \cap \operatorname{Ann}(\mathbf{A}) = 0 \right\},\,$$

which is stable under the action of Aut(A).

Now, let \mathbb{V} be an s-dimensional linear space and let us denote by $\mathbf{E}(\mathbf{A}, \mathbb{V})$ the set of all non-split s-dimensional central extensions of \mathbf{A} by \mathbb{V} . By above, we can write

$$\mathbf{E}\left(\mathbf{A}, \mathbb{V}\right) = \left\{\mathbf{A}_{\theta} : \theta\left(x, y\right) = \sum_{i=1}^{s} \theta_{i}\left(x, y\right) e_{i} \text{ and } \left\langle\left[\theta_{1}\right], \left[\theta_{2}\right], \dots, \left[\theta_{s}\right]\right\rangle \in \mathbf{T}_{s}(\mathbf{A})\right\}.$$

We also have the following result, which can be proved as in [28, Lemma 17].

Lemma 3. Let \mathbf{A}_{θ} , $\mathbf{A}_{\vartheta} \in \mathbf{E}(\mathbf{A}, \mathbb{V})$. Suppose that $\theta(x, y) = \sum_{i=1}^{s} \theta_{i}(x, y) e_{i}$ and $\vartheta(x, y) = \sum_{i=1}^{s} \vartheta_{i}(x, y) e_{i}$.

Then the left symmetric algebras A_{θ} and A_{ϑ} are isomorphic if and only if

$$\operatorname{Orb} \langle [\theta_1], [\theta_2], \dots, [\theta_s] \rangle = \operatorname{Orb} \langle [\vartheta_1], [\vartheta_2], \dots, [\vartheta_s] \rangle.$$

This shows that there exists a one-to-one correspondence between the set of $\operatorname{Aut}(\mathbf{A})$ -orbits on $\mathbf{T}_s(\mathbf{A})$ and the set of isomorphism classes of $\mathbf{E}(\mathbf{A}, \mathbb{V})$. Consequently we have a procedure that allows us, given a left symmetric algebra \mathbf{A}' of dimension n-s, to construct all non-split central extensions of \mathbf{A}' . This procedure is:

- (1) For a given left symmetric algebra \mathbf{A}' of dimension n-s, determine $\mathrm{H}^2(\mathbf{A}',\mathbb{C})$, $\mathrm{Ann}(\mathbf{A}')$ and $\mathrm{Aut}(\mathbf{A}')$.
- (2) Determine the set of Aut(A')-orbits on $T_s(A')$.

(3) For each orbit, construct the left symmetric algebra associated with a representative of it.

The above described method gives all (Novikov and non-Novikov) left symmetric algebras. But we are interested in developing this method in such a way that it only gives non-Novikov left symmetric algebras, because the classification of all Novikov algebras is given in [34]. Clearly, any central extension of a non-Novikov left symmetric algebra is non-Novikov. But a Novikov algebra may have extensions which are not Novikov algebras. More precisely, let N be a Novikov algebra and $\theta \in Z^2_L(N,\mathbb{C})$. Then N_θ is a Novikov algebra if and only if

$$\theta(xy,z) = \theta(xz,y).$$

for all $x, y, z \in \mathbb{N}$. Define the subspace $\mathrm{Z}^2_{\mathbb{N}}(\mathbb{N}, \mathbb{C})$ of $\mathrm{Z}^2_{\mathbb{L}}(\mathbb{N}, \mathbb{C})$ by

$$\mathbf{Z}^2_{\mathbf{N}}(\mathbf{N},\mathbb{C}) = \left\{ \ \theta \in \mathbf{Z}^2_{\mathbf{L}}(\mathbf{N},\mathbb{C}) : \theta(xy,z) = \theta(xz,y) \text{ for all } x,y,z \in \mathbf{N} \ \right\}.$$

Observe that $B^2(\mathbf{N},\mathbb{C})\subseteq Z_N^2(\mathbf{N},\mathbb{C})$. Let $H_N^2(\mathbf{N},\mathbb{C})=Z_N^2(\mathbf{N},\mathbb{C})\big/B^2(\mathbf{N},\mathbb{C})$. Then $H_N^2(\mathbf{N},\mathbb{C})$ is a subspace of $H_L^2(\mathbf{N},\mathbb{C})$. Define

$$\mathbf{R}_s(\mathbf{N}) = \left\{ \mathbf{W} \in \mathbf{T}_s(\mathbf{N}) : \mathbf{W} \in G_s(\mathrm{H}^2_\mathrm{N}(\mathbf{N}, \mathbb{C})) \right\},$$

$$\mathbf{U}_s(\mathbf{N}) = \left\{ \mathbf{W} \in \mathbf{T}_s(\mathbf{N}) : \mathbf{W} \notin G_s(\mathrm{H}^2_\mathrm{N}(\mathbf{N}, \mathbb{C})) \right\}.$$

Then $T_s(N) = R_s(N) \cup U_s(N)$. The sets $R_s(N)$ and $U_s(N)$ are stable under the action of $\operatorname{Aut}(N)$. Thus, the left symmetric algebras corresponding to the representatives of $\operatorname{Aut}(N)$ -orbits on $R_s(N)$ are Novikov algebras, while those corresponding to the representatives of $\operatorname{Aut}(N)$ -orbits on $U_s(N)$ are not Novikov algebras. Hence, we may construct all non-split non-Novikov left symmetric algebras A of dimension n with s-dimensional annihilator from a given left symmetric algebra A' of dimension n-s in the following way:

- (1) If A' is non-Novikov, then apply the Procedure.
- (2) Otherwise, do the following:
 - (a) Determine $U_s(A')$ and Aut(A').
 - (b) Determine the set of Aut(A')-orbits on $U_s(A')$.
 - (c) For each orbit, construct the left symmetric algebra corresponding to one of its representatives.
- 1.2. **Notations.** Let us introduce the following notations. Let \mathbf{A} be a nilpotent algebra with a basis e_1, e_2, \ldots, e_n . Then by Δ_{ij} we will denote the bilinear form $\Delta_{ij}: \mathbf{A} \times \mathbf{A} \longrightarrow \mathbb{C}$ with $\Delta_{ij}(e_l, e_m) = \delta_{il}\delta_{jm}$. The set $\{\Delta_{ij}: 1 \leq i, j \leq n\}$ is a basis for the linear space of bilinear forms on \mathbf{A} , so every $\theta \in \mathbf{Z}^2(\mathbf{A}, \mathbb{V})$ can be uniquely written as $\theta = \sum_{1 \leq i, j \leq n} c_{ij}\Delta_{ij}$, where $c_{ij} \in \mathbb{C}$. Let us fix the following notations:

$$\mathbf{L_{j}^{i}}$$
 — j th i -dimensional left symmetric (non-Novikov) algebra. $\mathbf{L_{i}^{i*}}$ — j th i -dimensional left symmetric (Novikov) algebra.

1.3. The algebraic classification of 3-dimensional nilpotent left symmetric algebras. There are no non-trivial 1-dimensional nilpotent left symmetric algebras. There is only one nontrivial 2-dimensional nilpotent left symmetric algebra (it is the non-split central extension of 1-dimensional algebra with zero product):

$$\mathbf{L}_{01}^{2*} : e_1 e_1 = e_2.$$

It is known the classification of all non-split 3-dimensional nilpotent left symmetric algebras:

$$\begin{array}{llll} \mathbf{L}_{02}^{3*} & : & e_1e_1=e_3 & e_2e_2=e_3 \\ \mathbf{L}_{03}^{3*} & : & e_1e_2=e_3 & e_2e_1=-e_3 \\ \mathbf{L}_{04}^{3*}(\lambda) & : & e_1e_1=\lambda e_3 & e_2e_1=e_3 & e_2e_2=e_3 \\ \mathbf{L}_{05}^{3*} & : & e_1e_1=e_2 & e_2e_1=e_3 \\ \mathbf{L}_{06}^{3*}(\lambda) & : & e_1e_1=e_2 & e_1e_2=e_3 & e_2e_1=\lambda e_3. \end{array}$$

1.4.1. The description of second cohomology spaces of 3-dimensional nilpotent left symmetric algebras. In the following table we give the description of the second cohomology space of 3-dimensional nilpotent left symmetric algebras.

A	$H^2_{f N}({f A})$	$\mathrm{H}^2_{\mathbf{L}}(\mathbf{A})$
\mathbf{L}_{01}^{3*}	$\left\langle [\Delta_{12}], [\Delta_{13}], [\Delta_{21}], [\Delta_{31}], [\Delta_{33}] \right\rangle$	$\left \operatorname{H}^2_{\mathbf{N}}(\mathbf{L}_{01}^{3*}) \oplus \left\langle [\Delta_{23}] \right angle \right $
\mathbf{L}_{02}^{3*}	$\left \left\langle [\Delta_{12}], [\Delta_{21}], [\Delta_{22}] \right\rangle \right $	$\left \operatorname{H}^2_{\mathbf{N}}(\mathbf{L}_{02}^{3*}) \oplus \left\langle [\Delta_{31}], [\Delta_{32}] \right\rangle$
\mathbf{L}_{03}^{3*}	$\left\langle [\Delta_{11}], [\Delta_{21}], [\Delta_{22}] \right angle$	$ H^2_{\mathbf{N}}(\mathbf{L}_{03}^{3*}) \oplus \left\langle [\Delta_{31} - 2\Delta_{13}], [\Delta_{32} - 2\Delta_{23}] \right\rangle $
$\mathbf{L}_{04}^{3*}(\lambda)_{\lambda \neq 0}$	$\Big \left\langle [\Delta_{11}], [\Delta_{12}], [\Delta_{21}] \right\rangle$	$H^2_{\mathbf{N}}(\mathbf{L}^{3*}_{04}(\lambda)) \oplus$
		$\left \left\langle [\Delta_{13}-\Delta_{31}-\Delta_{32}],[\Delta_{23}+\lambda\Delta_{31}]\right angle$
$\mathbf{L}_{04}^{3*}(0)$	$\left \left\langle [\Delta_{11}], [\Delta_{12}], [\Delta_{21}], [\Delta_{13} - \Delta_{31} - \Delta_{32}], [\Delta_{23}] \right\rangle \right $	$H^2_{\mathbf{N}}(\mathbf{L}_{04}^{3*}(0))$
${f L}_{05}^{3*}$	$\left\langle [\Delta_{12}], [\Delta_{13} - \Delta_{31}] \right\rangle$	$H^2_{\mathbf{N}}(\mathbf{L}_{05}^{3*}) \oplus \left\langle [\Delta_{22} + \Delta_{31}], [\Delta_{23}] \right\rangle$
$\mathbf{L}_{06}^{3*}(\lambda)$	$\left\langle [\Delta_{21}], [(2-\lambda)\Delta_{13} + \lambda(\Delta_{22} + \Delta_{31})] \right\rangle$	$\left \operatorname{H}^{2}_{\mathbf{N}}(\mathbf{L}_{06}^{3*}(\lambda)) \oplus \left\langle \left[\Delta_{22} + \Delta_{13} - \Delta_{31} \right] \right\rangle$

Remark 4. Since $H^2_{\mathbf{L}}(\mathbf{L}_{04}^{3*}(0)) = H^2_{\mathbf{N}}(\mathbf{L}_{04}^{3*}(0))$, then central extension of the algebra $\mathbf{L}_{04}^{3*}(0)$ give us only Novikov algebras.

1.4.2. Central extensions of L_{01}^{3*} . Let us use the following notations:

$$\nabla_1 = [\Delta_{12}], \quad \nabla_2 = [\Delta_{13}], \quad \nabla_3 = [\Delta_{21}], \quad \nabla_4 = [\Delta_{31}], \quad \nabla_5 = [\Delta_{33}], \quad \nabla_6 = [\Delta_{23}].$$

Take $\theta = \sum_{i=1}^{6} \alpha_i \nabla_i \in H^2_L(\mathbf{L}_{01}^{3*})$. The automorphism group of \mathbf{L}_{01}^{3*} consists of invertible matrices of the form

$$\phi = \begin{pmatrix} x & 0 & 0 \\ y & x^2 & u \\ z & 0 & t \end{pmatrix}.$$

Since

$$\phi^T \begin{pmatrix} 0 & \alpha_1 & \alpha_2 \\ \alpha_3 & 0 & \alpha_6 \\ \alpha_4 & 0 & \alpha_5 \end{pmatrix} \phi = \begin{pmatrix} \alpha^* & \alpha_1^* & \alpha_2^* \\ \alpha_3^* & 0 & \alpha_6^* \\ \alpha_4^* & 0 & \alpha_5^* \end{pmatrix},$$

we have that the action of $\operatorname{Aut}(\mathbf{L}_{01}^{3*})$ on the subspace $\langle \sum_{i=1}^{6} \alpha_i \nabla_i \rangle$ is given by $\langle \sum_{i=1}^{6} \alpha_i^* \nabla_i \rangle$, where

$$\begin{array}{llll} \alpha_1^* & = & \alpha_1 x^3 & & \alpha_2^* & = & \alpha_1 x u + \alpha_2 x t + \alpha_5 z t + \alpha_6 y t \\ \alpha_3^* & = & \alpha_3 x^3 + \alpha_6 x^2 z & & \alpha_4^* & = & \alpha_3 x u + \alpha_4 x t + \alpha_5 z t + \alpha_6 z u \\ \alpha_5^* & = & \alpha_5 t^2 + \alpha_6 t u & & \alpha_6^* & = & \alpha_6 x^2 t. \end{array}$$

Since $H^2_L(L^{3*}_{01}) = H^2_N(L^{3*}_{01}) \oplus \langle \nabla_6 \rangle$ and we are interested only in new algebras, we have $\alpha_6 \neq 0$. Then putting $z = -\frac{\alpha_3 x}{\alpha_6}$, $u = -\frac{\alpha_5 t}{\alpha_6}$ and $y = \frac{(\alpha_3 \alpha_5 + \alpha_1 \alpha_5 - \alpha_2 \alpha_6)x}{\alpha_6^2}$, we have

$$\alpha_2^* = \alpha_3^* = \alpha_5^* = 0$$
 $\alpha_1^* = \alpha_1 x^3$ $\alpha_4^* = \frac{(\alpha_4 \alpha_6 - \alpha_3 \alpha_5)xt}{\alpha_6}$ $\alpha_6^* = \alpha_6 x^2 t$.

Consider the following cases.

- (1) $\alpha_1 \neq 0$, $\alpha_4 \alpha_6 \alpha_3 \alpha_5 \neq 0$, then choosing $x = \frac{\alpha_4 \alpha_6 \alpha_3 \alpha_5}{\alpha_6^2}$, $t = \frac{\alpha_1 (\alpha_4 \alpha_6 \alpha_3 \alpha_5)}{\alpha_6^3}$, we have the representative $\langle \nabla_1 + \nabla_4 + \nabla_6 \rangle$.
- (2) $\alpha_1 = 0$, $\alpha_4 \alpha_6 \alpha_3 \alpha_5 \neq 0$, then choosing $x = \frac{\alpha_4 \alpha_6 \alpha_3 \alpha_5}{\alpha_6^2}$, we have the representative $\langle \nabla_4 + \nabla_6 \rangle$.
- (3) $\alpha_1 \neq 0, \alpha_4 \alpha_6 \alpha_3 \alpha_5 = 0$, then choosing $t = \frac{\alpha_1 x}{\alpha_6}$, we have the representative $\langle \nabla_1 + \nabla_6 \rangle$.
- (4) $\alpha_1 = 0$, $\alpha_4 \alpha_6 \alpha_3 \alpha_5 = 0$, then we have the representative $\langle \nabla_6 \rangle$.

Hence, we have the following distints orbits

$$\langle \nabla_1 + \nabla_4 + \nabla_6 \rangle \quad \langle \nabla_4 + \nabla_6 \rangle \quad \langle \nabla_1 + \nabla_6 \rangle \quad \langle \nabla_6 \rangle$$

which give the following new algebras:

1.4.3. Central extensions of L_{02}^{3*} . Let us use the following notations:

$$\nabla_1 = [\Delta_{12}], \quad \nabla_2 = [\Delta_{21}], \quad \nabla_3 = [\Delta_{22}], \quad \nabla_4 = [\Delta_{31}], \quad \nabla_5 = [\Delta_{32}].$$

Take $\theta = \sum_{i=1}^{5} \alpha_i \nabla_i \in H^2_L(\mathbf{L}_{02}^{3*})$. The automorphism group of \mathbf{L}_{02}^{3*} consists of invertible matrices of the form

$$\phi_1 = \begin{pmatrix} x & -y & 0 \\ y & x & 0 \\ z & t & x^2 + y^2 \end{pmatrix} \quad \text{or} \quad \phi_2 = \begin{pmatrix} x & y & 0 \\ y & -x & 0 \\ z & t & x^2 + y^2 \end{pmatrix}.$$

Since

$$\phi_1^T \begin{pmatrix} 0 & \alpha_1 & 0 \\ \alpha_2 & \alpha_3 & 0 \\ \alpha_4 & \alpha_5 & 0 \end{pmatrix} \phi_1 = \begin{pmatrix} \alpha^* & \alpha_1^* & 0 \\ \alpha_2^* & \alpha^* + \alpha_3^* & 0 \\ \alpha_4^* & \alpha_5^* & 0 \end{pmatrix},$$

we have that the action of $\operatorname{Aut}(\mathbf{L}_{02}^{3*})^+$ (it is the subgroup in $\operatorname{Aut}(\mathbf{L}_{02}^{3*})$ formed by all automorphisms of the first type)on the subspace $\langle \sum_{i=1}^{5} \alpha_i \nabla_i \rangle$ is given by $\langle \sum_{i=1}^{5} \alpha_i^* \nabla_i \rangle$, where

$$\alpha_1^* = \alpha_1 x^2 - \alpha_2 y^2 + \alpha_3 xy - \alpha_4 yz + \alpha_5 xz
\alpha_2^* = -\alpha_1 y^2 + \alpha_2 x^2 + \alpha_3 xy + \alpha_4 xt + \alpha_5 yt
\alpha_3^* = -2\alpha_1 xy - 2\alpha_2 xy + \alpha_3 (x^2 - y^2) - \alpha_4 (xz + yt) - \alpha_5 (yz - xt)
\alpha_4^* = (\alpha_4 x + \alpha_5 y) (x^2 + y^2)
\alpha_5^* = (-\alpha_4 y + \alpha_5 x) (x^2 + y^2).$$

Since $H^2_{\mathbf{L}}(\mathbf{L}_{02}^{3*})=H^2_{\mathbf{N}}(\mathbf{L}_{02}^{3*})\oplus\langle\nabla_4,\nabla_5\rangle$ and we are interested only in new algebras, we have $(\alpha_4,\alpha_5)\neq$

(0,0). Moreover, without loss of generality, one can assume $\alpha_4 \neq 0$. Then we have the following cases.

(1) $\alpha_4^2 + \alpha_5^2 \neq 0$, then choosing $y = \frac{x\alpha_5}{\alpha_4}$, $t = \frac{(\alpha_1\alpha_5^2 - \alpha_2\alpha_4^2 - \alpha_3\alpha_4\alpha_5)x}{\alpha_4(\alpha_4^2 + \alpha_5^2)}$, $z = \frac{(\alpha_3(\alpha_4^2 - \alpha_5^2) - 2\alpha_4\alpha_5(\alpha_1 + \alpha_2))x}{\alpha_4(\alpha_4^2 + \alpha_5^2)}$, we have

$$\alpha_2^* = \alpha_3^* = \alpha_5^* = 0 \quad \alpha_1^* = \frac{(\alpha_1 \alpha_4^2 - \alpha_2 \alpha_5^2 + \alpha_3 \alpha_4 \alpha_5) x^2}{\alpha_4^2} \quad \alpha_4^* = \frac{x^3 (\alpha_4^2 + \alpha_5^2)^2}{\alpha_4^3}.$$

- (a) if $\alpha_1\alpha_4^2 \alpha_2\alpha_5^2 + \alpha_3\alpha_4\alpha_5 = 0$, then we have the representative $\langle \nabla_4 \rangle$; (b) if $\alpha_1\alpha_4^2 \alpha_2\alpha_5^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, then choosing $x = \frac{\alpha_4(\alpha_1\alpha_4^2 \alpha_2\alpha_5^2 + \alpha_3\alpha_4\alpha_5)}{(\alpha_4^2 + \alpha_5^2)^2}$, we have the representative $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, then choosing $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, then choosing $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 = 0$, we have the representative $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, where $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, we have the representative $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, where $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, we have the representative $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, where $\alpha_1\alpha_2^2 \alpha_3\alpha_4\alpha_5 \neq 0$, where $\alpha_1\alpha_2^2 \alpha_2\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, where $\alpha_1\alpha_2^2 \alpha_3\alpha_3^2 + \alpha_3\alpha_4\alpha_5 \neq 0$, where $\alpha_1\alpha_2^2 \alpha_3\alpha_4\alpha_5$ tative $\langle \nabla_1 + \nabla_4 \rangle$;
- (2) $\alpha_4^2 + \alpha_5^2 = 0$, i.e., $\alpha_5 = \pm i\alpha_4$, then choosing

$$t = \frac{\alpha_1 y^2 - \alpha_2 x^2 - \alpha_3 xy}{\alpha_4 (x \pm iy)}, \ z = \frac{-2\alpha_1 xy - 2\alpha_2 xy + \alpha_3 (x^2 - y^2) + \alpha_4 (-y \pm ix)}{\alpha_4 (x \pm iy)},$$

we have

$$\alpha_2^* = \alpha_3^* = 0, \quad \alpha_1^* = -\frac{(\alpha_1 + \alpha_2 \pm i\alpha_3)(x^2 + y^2)^2}{(x \pm iy)^2},$$

$$\alpha_4^* = (x^2 + y^2)(x \pm iy)\alpha_4, \quad \alpha_5^* = \pm i(x^2 + y^2)(x \pm iy)\alpha_4.$$

- (a) $\alpha_1 + \alpha_2 \pm i\alpha_3 = 0$, then we have representative $\langle \nabla_4 \pm i\nabla_5 \rangle$. (b) $\alpha_1 + \alpha_2 \pm i\alpha_3 \neq 0$, then choosing $x = -\frac{\alpha_1 + \alpha_2 \pm i\alpha_3}{\alpha_4}$ and y = 0, we have representative $\langle \nabla_1 + \nabla_4 \pm i \nabla_5 \rangle$.

Since the automorphism $\phi = diag(1, -1, 1)$ acts as

$$\phi(\nabla_4 + i\nabla_5) = \langle \nabla_4 - i\nabla_5 \rangle$$
 and $\phi(\nabla_1 + \nabla_4 + i\nabla_5) = \langle \nabla_1 + \nabla_4 - i\nabla_5 \rangle$,

we have two representatives of distinct orbits $\langle \nabla_4 + i \nabla_5 \rangle$ and $\langle \nabla_1 + \nabla_4 + i \nabla_5 \rangle$. Hence, we have the following distints orbits

$$\langle \nabla_4 \rangle \quad \langle \nabla_1 + \nabla_4 \rangle \quad \langle \nabla_4 + i \nabla_5 \rangle \quad \langle \nabla_1 + \nabla_4 + i \nabla_5 \rangle,$$

which give the following new algebras:

1.4.4. Central extensions of L_{03}^{3*} . Let us use the following notations:

$$\nabla_1 = [\Delta_{11}], \quad \nabla_2 = [\Delta_{21}], \quad \nabla_3 = [\Delta_{22}], \quad \nabla_4 = [\Delta_{31} - 2\Delta_{13}], \quad \nabla_5 = [\Delta_{32} - 2\Delta_{23}].$$

Take $\theta = \sum_{i=0}^{5} \alpha_i \nabla_i \in H^2_L(\mathbf{L}_{03}^{3*})$. The automorphism group of \mathbf{L}_{03}^{3*} consists of invertible matrices of the form

$$\phi = \begin{pmatrix} x & u & 0 \\ y & v & 0 \\ z & t & xv - yu \end{pmatrix}.$$

Since

$$\phi^{T} \begin{pmatrix} \alpha_{1} & 0 & -2\alpha_{4} \\ \alpha_{2} & \alpha_{3} & -2\alpha_{5} \\ \alpha_{4} & \alpha_{5} & 0 \end{pmatrix} \phi = \begin{pmatrix} \alpha_{1}^{*} & \alpha^{*} & -2\alpha_{4}^{*} \\ \alpha_{2}^{*} - \alpha^{*} & \alpha_{3}^{*} & -2\alpha_{5}^{*} \\ \alpha_{4}^{*} & \alpha_{5}^{*} & 0 \end{pmatrix},$$

we have that the action of $\operatorname{Aut}(\mathbf{L}_{03}^{3*})$ on the subspace $\langle \sum_{i=1}^{3} \alpha_i \nabla_i \rangle$ is given by $\langle \sum_{i=1}^{3} \alpha_i^* \nabla_i \rangle$, where

$$\alpha_1^* = \alpha_1 x^2 + \alpha_2 xy + \alpha_3 y^2 - \alpha_4 xz - \alpha_5 yz
\alpha_2^* = 2\alpha_1 xu + \alpha_2 (xv + yu) + 2\alpha_3 yv - \alpha_4 (xt + zu) - \alpha_5 (yt + zv)
\alpha_3^* = \alpha_1 u^2 + \alpha_2 uv + \alpha_3 v^2 - \alpha_4 ut - \alpha_5 vt
\alpha_4^* = (\alpha_4 x + \alpha_5 y) (xv - yu)
\alpha_5^* = (\alpha_4 u + \alpha_5 v) (xv - yu).$$

Since $\mathrm{H}^2_{\mathbf{L}}(\mathbf{L}^{3*}_{03})=\mathrm{H}^2_{\mathbf{N}}(\mathbf{L}^{3*}_{03})\oplus\langle\nabla_4,\nabla_5\rangle$, we have $(\alpha_4,\alpha_5)\neq(0,0)$. Moreover, without loss of generality, one can assume $\alpha_4\neq0$. Choosing $u=-\frac{\alpha_5v}{\alpha_4},\ z=\frac{\alpha_1x^2+\alpha_2xy+\alpha_3y^2}{\alpha_4x+\alpha_5y}$ and $t=\frac{2\alpha_1xu+\alpha_2(xv+yu)+2\alpha_3yv}{\alpha_4x+\alpha_5y}$, we obtain

$$\alpha_1^* = \alpha_2^* = \alpha_5^* = 0, \quad \alpha_3^* = \frac{(\alpha_1 \alpha_5^2 - \alpha_2 \alpha_4 \alpha_5 + \alpha_3 \alpha_4^2)v^2}{\alpha_4^2}, \quad \alpha_4^* = \frac{v(\alpha_4 x + \alpha_5 y)^2}{\alpha_4}.$$

Then we have the following cases.

- (1) $\alpha_1\alpha_5^2 \alpha_2\alpha_4\alpha_5 + \alpha_3\alpha_4^2 = 0$, then we have the representative $\langle \nabla_4 \rangle$. (2) $\alpha_1\alpha_5^2 \alpha_2\alpha_4\alpha_5 + \alpha_3\alpha_4^2 \neq 0$, then choosing $v = \frac{\alpha_4(\alpha_4x + \alpha_5y)^2}{(\alpha_1\alpha_5^2 \alpha_2\alpha_4\alpha_5 + \alpha_3\alpha_4^2)}$, we have the representative $\langle \nabla_3 + \nabla_4 \rangle$.

Summarizing, we have the following distinct orbits

$$\langle \nabla_4 \rangle \quad \langle \nabla_3 + \nabla_4 \rangle.$$

Hence, we have the following new algebras:

1.4.5. Central extensions of $L_{04}^{3*}(\lambda)_{\lambda\neq 0}$. Let us use the following notations:

$$\nabla_1 = [\Delta_{12}], \quad \nabla_2 = [\Delta_{21}], \quad \nabla_3 = [\Delta_{22}], \quad \nabla_4 = [\Delta_{13} - \Delta_{31} - \Delta_{32}], \quad \nabla_5 = [\Delta_{23} + \lambda \Delta_{31}].$$

Take $\theta = \sum_{i=1}^{5} \alpha_i \nabla_i \in \mathrm{H}^2_{\mathbf{L}}(\mathbf{L}_{04}^{3*}(\lambda)_{\lambda \neq 0})$. The automorphism group of $\mathbf{L}_{04}^{3*}(\lambda)_{\lambda \neq 0}$ consists of invertible matrices of the form

$$\phi = \begin{pmatrix} x & y & 0 \\ -\lambda y & x - y & 0 \\ z & t & x^2 - xy + \lambda y^2 \end{pmatrix}.$$

Since

$$\phi^{T} \begin{pmatrix} 0 & \alpha_{1} & \alpha_{4} \\ \alpha_{2} & \alpha_{3} & \alpha_{5} \\ -\alpha_{4} + \lambda \alpha_{5} & -\alpha_{4} & 0 \end{pmatrix} \phi = \begin{pmatrix} \lambda \alpha^{*} & \alpha_{1}^{*} & \alpha_{4}^{*} \\ \alpha^{*} + \alpha_{2}^{*} & \alpha^{*} + \alpha_{3}^{*} & \alpha_{5}^{*} \\ -\alpha_{4}^{*} + \lambda \alpha_{5}^{*} & -\alpha_{4}^{*} & 0 \end{pmatrix},$$

we have that the action of $\operatorname{Aut}(\mathbf{L}_{04}^{3*}(\lambda)_{\lambda\neq 0})$ on the subspace $\langle \sum_{i=1}^{5} \alpha_i \nabla_i \rangle$ is given by $\langle \sum_{i=1}^{5} \alpha_i^* \nabla_i \rangle$, where

$$\begin{array}{lll} \alpha_1^* & = & \alpha_1(x^2-xy) - \alpha_2\lambda y^2 - \alpha_3(\lambda xy - \lambda y^2) + \alpha_4(xt-xz) + \alpha_5(\lambda yz - \lambda yt), \\ \alpha_2^* & = & \alpha_1(xy-\lambda y^2) + \alpha_2x^2 - \alpha_3\lambda xy + \alpha_4(\lambda yt-xt) + \alpha_5\lambda xt, \\ \alpha_3^* & = & \alpha_1(2xy-y^2) + \alpha_2(2xy-y^2) + \alpha_3\left((x-y)^2 - \lambda y^2\right) \\ & & + \alpha_4(yt-xt-yz) + \alpha_5(\lambda yt+xt-yt-xz+yz), \\ \alpha_4^* & = & (\alpha_4x-\alpha_5\lambda y)(x^2-xy+\lambda y^2), \\ \alpha_5^* & = & (\alpha_4y+\alpha_5(x-y))(x^2-xy+\lambda y^2). \end{array}$$

Since $H^2_{\mathbf{L}}(\mathbf{L}_{04}^{3*}(\lambda)_{\lambda\neq 0}) = H^2_{\mathbf{N}}(\mathbf{L}_{04}^{3*}(\lambda)_{\lambda\neq 0}) \oplus \langle \nabla_4, \nabla_5 \rangle$, we have $(\alpha_4, \alpha_5) \neq (0, 0)$. Then we have the following cases.

(1)
$$\alpha_4 = 0$$
, then $\alpha_5 \neq 0$ and choosing $y = 0$, $t = -\frac{\alpha_2 x}{\lambda \alpha_5}$, $z = \frac{(\alpha_3 \lambda - \alpha_2)x}{\alpha_5 \lambda}$, we have

$$\alpha_2^* = \alpha_3^* = \alpha_4^* = 0, \quad \alpha_1^* = \alpha_1 x^2, \quad \alpha_5^* = \alpha_5 x^3.$$

- (a) $\alpha_1 = 0$, then we have the representative $\langle \nabla_5 \rangle$.
- (b) $\alpha_1 \neq 0$, then choosing $x = \frac{\alpha_1}{\alpha_5}$, we have the representative $\langle \nabla_1 + \nabla_5 \rangle$.
- (2) $\alpha_4 \neq 0$, and $\alpha_4^2 \alpha_4 \alpha_5 + \alpha_5^2 \lambda \neq 0$, then choosing $x = \frac{\lambda \alpha_5}{\alpha_4}$ and y = 1, we have $\alpha_4^* = 0$ and it is the situation considered above.
- (3) $\alpha_4 \neq 0$ and $\alpha_4^2 \alpha_4 \alpha_5 + \alpha_5^2 \lambda = 0$, then choosing

$$y = 0$$
 $t = \frac{\alpha_2 x}{\alpha_4 - \lambda \alpha_5}$ $z = \frac{(\alpha_1 \alpha_4 - \lambda \alpha_1 \alpha_5 + \alpha_2 \alpha_4)x}{\alpha_4 (\alpha_4 - \lambda \alpha_5)}$

we have

$$\alpha_1^* = \alpha_2^* = 0 \quad \alpha_3^* = \frac{\left((\alpha_4 - \lambda \alpha_5)(\alpha_3 \alpha_4 - \alpha_1 \alpha_5) - \alpha_2 \alpha_4^2\right) x^2}{\alpha_4(\alpha_4 - \lambda \alpha_5)} \quad \alpha_4^* = x^3 \alpha_4 \quad \alpha_5^* = x^3 \alpha_5.$$

- (a) $\alpha_2 \alpha_4^2 = (\alpha_4 \lambda \alpha_5)(\alpha_3 \alpha_4 \alpha_1 \alpha_5)$, then we have the representative $\langle \nabla_4 + \frac{1 \pm \sqrt{1 4\lambda}}{2\lambda} \nabla_5 \rangle$.
- (b) $\alpha_2 \alpha_4^2 \neq (\alpha_4 \lambda \alpha_5)(\alpha_3 \alpha_4 \alpha_1 \alpha_5)$, then we have the representative $\langle \frac{1}{2\lambda} \nabla_3 + \nabla_4 + \frac{1 \pm \sqrt{1 4\lambda}}{2\lambda} \nabla_5 \rangle$.

Summarizing, we have the following distinct orbits

$$\langle \nabla_5 \rangle \qquad \langle 2\lambda \nabla_4 + (1 - \sqrt{1 - 4\lambda}) \nabla_5 \rangle \qquad \langle \nabla_3 + 2\lambda \nabla_4 + (1 - \sqrt{1 - 4\lambda}) \nabla_5 \rangle$$

$$\langle \nabla_1 + \nabla_5 \rangle \qquad \langle 2\lambda \nabla_4 + (1 + \sqrt{1 - 4\lambda}) \nabla_5 \rangle \qquad \langle \nabla_3 + 2\lambda \nabla_4 + (1 + \sqrt{1 - 4\lambda}) \nabla_5 \rangle.$$

Hence, we have the following new algebras:

1.4.6. Central extensions of L_{05}^{3*} . Let us use the following notations:

$$\nabla_1 = [\Delta_{12}], \quad \nabla_2 = [\Delta_{13} - \Delta_{31}], \quad \nabla_3 = [\Delta_{22} + \Delta_{31}], \quad \nabla_4 = [\Delta_{23}].$$

Take $\theta = \sum_{i=1}^{4} \alpha_i \nabla_i \in H^2_L(L^{3*}_{05})$. The automorphism group of L^{3*}_{05} consists of invertible matrices of the form

$$\phi = \begin{pmatrix} x & 0 & 0 \\ y & x^2 & 0 \\ z & xy & x^3 \end{pmatrix}.$$

Since

$$\phi^{T} \begin{pmatrix} 0 & \alpha_{1} & \alpha_{2} \\ 0 & \alpha_{3} & \alpha_{4} \\ -\alpha_{2} + \alpha_{3} & 0 & 0 \end{pmatrix} \phi = \begin{pmatrix} \alpha^{*} & \alpha_{1}^{*} & \alpha_{2}^{*} \\ \alpha^{**} & \alpha_{3}^{*} & \alpha_{4}^{*} \\ -\alpha_{2}^{*} + \alpha_{3}^{*} & 0 & 0 \end{pmatrix},$$

we have that the action of $\operatorname{Aut}(\mathbf{L}_{05}^{3*})$ on the subspace $\langle \sum_{i=1}^{4} \alpha_i \nabla_i \rangle$ is given by $\langle \sum_{i=1}^{4} \alpha_i^* \nabla_i \rangle$, where

$$\begin{array}{ll} \alpha_1^* = \alpha_1 x^3 + (\alpha_2 + \alpha_3) x^2 y + \alpha_4 x y^2 & \alpha_2^* = \alpha_2 x^4 + \alpha_4 x^3 y \\ \alpha_3^* = \alpha_3 x^4 + \alpha_4 x^3 y & \alpha_4^* = \alpha_4 x^5. \end{array}$$

Since $H^2_L(L^{3*}_{05}) = H^2_N(L^{3*}_{05}) \oplus \langle \nabla_3, \nabla_4 \rangle$, we have $(\alpha_3, \alpha_4) \neq (0, 0)$. Then we have the following cases.

- (1) $\alpha_4 = 0$, then
 - (a) if $\alpha_2 + \alpha_3 \neq 0$, then choosing $y = -\frac{\alpha_1 x}{\alpha_2 + \alpha_3}$, we have the representative $\langle \alpha \nabla_2 + \nabla_3 \rangle_{\alpha \neq -1}$;
 - (b) if $\alpha_2 + \alpha_3 = 0$, then we have the representatives $\langle -\nabla_2 + \nabla_3 \rangle$ and $\langle \nabla_1 \nabla_2 + \nabla_3 \rangle$ depending on whether $\alpha_1 = 0$ or not.
- (2) if $\alpha_4 \neq 0$, then choosing $y = -\frac{\alpha_3 x}{\alpha_4}$, we have

$$\alpha_1^* = \frac{(\alpha_1 \alpha_4 - \alpha_2 \alpha_3)x^3}{\alpha_4}, \quad \alpha_2^* = (\alpha_2 - \alpha_3)x^4, \quad \alpha_3^* = 0, \quad \alpha_4^* = \alpha_4 x^5,$$

- (a) if $\alpha_1 \alpha_4 \alpha_2 \alpha_3 = 0$, $\alpha_2 \alpha_3 = 0$, then we have the representative $\langle \nabla_4 \rangle$;
- (b) if $\alpha_1\alpha_4 \alpha_2\alpha_3 = 0$, $\alpha_2 \alpha_3 \neq 0$, then we have the representative $\langle \nabla_2 + \nabla_4 \rangle$;
- (c) if $\alpha_1\alpha_4 \alpha_2\alpha_3 \neq 0$, then we have the representative $\langle \nabla_1 + \alpha \nabla_2 + \nabla_4 \rangle$.

Hence, we have the following distints orbits

$$\langle \alpha \nabla_2 + \nabla_3 \rangle \quad \langle \nabla_1 - \nabla_2 + \nabla_3 \rangle \quad \langle \nabla_4 \rangle \quad \langle \nabla_2 + \nabla_4 \rangle \quad \langle \nabla_1 + \alpha \nabla_2 + \nabla_4 \rangle,$$

which give the following new algebras:

1.4.7. Central extensions of $L_{06}^{3*}(\lambda)$. Let us use the following notations:

$$\nabla_1 = [\Delta_{21}] \quad \nabla_2 = [(2 - \lambda)\Delta_{13} + \lambda\Delta_{22} + \lambda\Delta_{31}] \quad \nabla_3 = [\Delta_{22} + \Delta_{13} - \Delta_{31}].$$

Take $\theta = \sum_{i=1}^{3} \alpha_i \nabla_i \in \mathrm{H}^2_{\mathbf{L}}(\mathbf{L}_{06}^{3*}(\lambda))$. The automorphism group of $\mathbf{L}_{06}^{3*}(\lambda)$ consists of invertible matrices of the form

$$\phi = \begin{pmatrix} x & 0 & 0 \\ y & x^2 & 0 \\ z & xy(1+\lambda) & x^3 \end{pmatrix}.$$

Since

$$\phi^{T} \begin{pmatrix} 0 & 0 & (2-\lambda)\alpha_{2} + \alpha_{3} \\ \alpha_{1} & \lambda\alpha_{2} + \alpha_{3} & 0 \\ \lambda\alpha_{2} - \alpha_{3} & 0 & 0 \end{pmatrix} \phi = \begin{pmatrix} \alpha^{**} & \alpha^{*} & (2-\lambda)\alpha_{2}^{*} + \alpha_{3}^{*} \\ \alpha_{1}^{*} + \lambda\alpha^{*} & \lambda\alpha_{2}^{*} + \alpha_{3}^{*} & 0 \\ \lambda\alpha_{2}^{*} - \alpha_{3}^{*} & 0 & 0 \end{pmatrix},$$

we have that the action of $\operatorname{Aut}(\mathbf{L}_{06}^{3*}(\lambda))$ on the subspace $\langle \sum_{i=1}^{3} \alpha_i \nabla_i \rangle$ is given by $\langle \sum_{i=1}^{3} \alpha_i^* \nabla_i \rangle$, where

$$\alpha_1^* = \alpha_1 x^3 + (\alpha_2(\lambda^3 - \lambda^2) - \alpha_3(\lambda^2 + 3\lambda))x^2y \quad \alpha_2^* = \alpha_2 x^4 \quad \alpha_3^* = \alpha_3 x^4.$$

Since $H^2_L(L^{3*}_{06}(\lambda)) = H^2_N(L^{3*}_{06}(\lambda)) \oplus \langle \nabla_3 \rangle$, we have $\alpha_3 \neq 0$. Then we have the following cases

- (1) if $\lambda=0$, then we have the representatives $\langle \alpha \nabla_2 + \nabla_3 \rangle$ and $\langle \nabla_1 + \alpha \nabla_2 + \nabla_3 \rangle$ depending on whether $\alpha_1=0$ or not.
- (2) if $\lambda = 1$, then choosing $y = \frac{\alpha_1 x}{4\alpha_3}$, we have the representative $\langle \alpha \nabla_2 + \nabla_3 \rangle$.
- (3) if $\lambda \notin \{0; 1\}$, then:
 - (a) if $\alpha_2(\lambda^3 \lambda^2) \alpha_3(\lambda^2 + 3\lambda) \neq 0$ then choosing $y = \frac{\alpha_1 x}{\alpha_3(\lambda^2 + 3\lambda) \alpha_2(\lambda^3 \lambda_2)}$, we have the representative $\langle \alpha \nabla_2 + \nabla_3 \rangle_{\alpha \neq \frac{\lambda + 3}{\lambda(\lambda 1)}}$;
 - (b) if $\alpha_2(\lambda^3 \lambda^2) \alpha_3(\lambda^2 + 3\lambda) = 0$ then we have the representative $\langle \frac{\lambda+3}{\lambda(\lambda-1)} \nabla_2 + \nabla_3 \rangle$ and $\langle \frac{1}{\lambda(\lambda-1)} \nabla_1 + \frac{\lambda+3}{\lambda(\lambda-1)} \nabla_2 + \nabla_3 \rangle$ depending on whether $\alpha_1 = 0$ or not.

Thus, we have the following orbits:

$$\langle \nabla_1 + \alpha \nabla_2 + \nabla_3 \rangle_{\lambda=0} \quad \langle \alpha \nabla_2 + \nabla_3 \rangle \quad \langle \nabla_1 + (\lambda+3) \nabla_2 + \lambda(\lambda-1) \nabla_3 \rangle_{\lambda \neq 0;1}.$$

Hence, we have the following new algebras:

1.5. **Classification theorem.** Now we are ready summarize all results related to the algebraic classification of complex 4-dimensional nilpotent left symmetric algebras.

Theorem A. Let \mathbf{L} be a complex 4-dimensional nilpotent left symmetric algebra. Then \mathbf{L} is a Novikov algebra or isomorphic to one algebra from the following list:

\mathbf{L}^4_{01}	:	$e_1e_1=e_2$	$e_1e_2=e_4$	$e_2e_3=e_4$	$e_3e_1=e_4$	
\mathbf{L}_{02}^4	:	$e_1e_1 = e_2$	$e_2e_3=e_4$	$e_3e_1=e_4$		
\mathbf{L}_{03}^{4}	:	$e_1e_1 = e_2$	$e_1e_2 = e_4$	$e_2e_3=e_4$		
\mathbf{L}_{04}^4	:	$e_1e_1 = e_2$	$e_2e_3 = e_4$			
\mathbf{L}_{05}^4	:	$e_1e_1 = e_3$	$e_2e_2=e_3$	$e_3e_1=e_4$		
\mathbf{L}_{06}^4	:	$e_1e_1 = e_3$	$e_1e_2 = e_4$	$e_2e_2=e_3$	$e_3e_1=e_4$	
\mathbf{L}_{07}^4	:	$e_1e_1 = e_3$	$e_2e_2=e_3$	$e_3e_1=e_4$	$e_3e_2 = ie_4$	
\mathbf{L}_{08}^4	:	$e_1e_1 = e_3$	$e_1e_2 = e_4$	$e_2e_2=e_3$	$e_3e_1 = e_4$	$e_3e_2 = ie_4$
\mathbf{L}_{09}^4	:	$e_1e_2 = e_3$	$e_1e_3 = -2e_4$	$e_2e_1 = -e_3$	$e_3e_1=e_4$	
\mathbf{L}_{10}^4	:	$e_1e_2 = e_3$	$e_1e_3 = -2e_4$	$e_2e_1 = -e_3$	$e_2e_2=e_4$	$e_3e_1=e_4$
$\mathbf{L}_{11}^4(\lambda)_{\lambda\neq 0}$:	$e_1e_1 = \lambda e_3$	$e_2e_1=e_3$	$e_2e_2=e_3$	$e_2e_3=e_4$	$e_3 e_1 = \lambda e_4$
$\mathbf{L}_{12}^4(\lambda)_{\lambda \neq 0}$:	$e_1e_1 = \lambda e_3$	$e_1e_2 = e_4$	$e_2e_1=e_3$		
		$e_2e_2=e_3$	$e_2e_3=e_4$	$e_3e_1 = \lambda e_4$		
$\mathbf{L}_{13}^4(\lambda)_{\lambda \neq 0}$:	$e_1e_1 = \lambda e_3$	$e_2e_2=e_3$	$e_2e_3 = (1 - \sqrt{1 - 1})^2$	$\sqrt{1-4\lambda})e_4$	$e_3 e_2 = -2\lambda e_4$
		$e_2e_1=e_3$	$e_1 e_3 = 2\lambda e_4$	$e_3e_1=-\lambda(1-$	$+\sqrt{1-4\lambda})e_4$	
$\mathbf{L}_{14}^4(\lambda)_{\lambda \neq 0}$:	$e_1e_1 = \lambda e_3$	$e_2e_2 = e_3 + e_4$	$e_2e_3 = (1 - \sqrt{1 - v})$	$\sqrt{1-4\lambda})e_4$	$e_3e_2 = -2\lambda e_4$
		$e_2e_1=e_3$	$e_1e_3 = 2\lambda e_4$	$e_3e_1=-\lambda(1-$	$+\sqrt{1-4\lambda})e_4$	
$\mathbf{L}_{15}^4(\lambda)_{\lambda\neq 0}$:	$e_1e_1 = \lambda e_3$	$e_2e_2=e_3$	$e_2e_3 = (1 + \sqrt{1 + \sqrt{1 + v^2}})$	$\sqrt{1-4\lambda})e_4$	$e_3 e_2 = -2\lambda e_4$
		$e_2e_1=e_3$	$e_1 e_3 = 2\lambda e_4$	$e_3e_1 = -\lambda(1 - \epsilon_3)$	$-\sqrt{1-4\lambda}e_4$	
$\mathbf{L}_{16}^4(\lambda)_{\lambda\neq 0}$:	$e_1e_1 = \lambda e_3$	$e_2e_2 = e_3 + e_4$	$e_2e_3 = (1 + \sqrt{1 + v})^2$	$\sqrt{1-4\lambda})e_4$	$e_3e_2 = -2\lambda e_4$
		$e_2e_1=e_3$	$e_1e_3 = 2\lambda e_4$	$e_3e_1=-\lambda(1-$	$-\sqrt{1-4\lambda}e_4$	
$\mathbf{L}_{17}^4(\alpha)$:	$e_1e_1 = e_2$	$e_1e_3 = \alpha e_4$	$e_2e_1=e_3$	$e_2e_2=e_4$	$e_3e_1 = (1 - \alpha)e_4$
\mathbf{L}_{18}^{4}	:	$e_1e_1=e_2$	$e_1e_2=e_4$	$e_1e_3 = -e_4$		
		$e_2e_1=e_3$	$e_2e_2=e_4$	$e_3e_1=2e_4$		
\mathbf{L}_{19}^4	:	$e_1e_1=e_2$	$e_2e_1=e_3$	$e_2e_3=e_4$		
$\frac{\overline{\mathbf{L}_{20}^{\frac{13}{4}}}}{\mathbf{L}_{21}^{4}(\alpha)}$:	$e_1e_1 = e_2$	$e_1e_3=e_4$	$e_2e_1=e_3$	$e_2e_3=e_4$	$e_3e_1 = -e_4$
$\mathbf{L}_{21}^4(\alpha)$:	$e_1e_1=e_2$	$e_1e_2=e_4$	$e_1e_3 = \alpha e_4$		
		$e_2e_1=e_3$	$e_2e_3=e_4$	$e_3e_1 = -\alpha e_4$		
$\mathbf{L}_{22}^4(\alpha)$:	$e_1e_1=e_2$	$e_1e_2=e_3$	$e_1 e_3 = (2\alpha +$	$1)e_4$	
		$e_2e_1=e_4$	$e_2e_2=e_4$	$e_3e_1 = -e_4$		
$\mathbf{L}_{23}^4(\lambda,\alpha)$:	$e_1e_1 = e_2$	$e_1e_2 = e_3$	$e_1e_3 = ((2 - \lambda e_1)^2)^2$,	
- 4()		$e_2e_1 = \lambda e_3$	$e_2 e_2 = (\lambda \alpha + 1) e_4$	$e_3e_1 = (\lambda \alpha -$		
$\mathbf{L}_{24}^4(\lambda)_{\lambda \notin \{0;1\}}$:	$e_1e_1 = e_2$	$e_1e_2 = e_3$	$e_1e_3 = 2(3 - 1)$	$(\lambda)e_4$	
		$e_2e_1 = \lambda e_3 + e_4$	$e_2 e_2 = 2\lambda(\lambda + 1)e_4$	$e_3e_1 = 4\lambda e_4$		

2. THE GEOMETRIC CLASSIFICATION OF NILPOTENT LEFT SYMMETRIC ALGEBRAS

2.1. **Definitions and notation.** Given an n-dimensional vector space \mathbb{V} , the set $\operatorname{Hom}(\mathbb{V} \otimes \mathbb{V}, \mathbb{V}) \cong \mathbb{V}^* \otimes \mathbb{V}^* \otimes \mathbb{V}$ is a vector space of dimension n^3 . This space has the structure of the affine variety \mathbb{C}^{n^3} . Indeed, let us fix a basis e_1, \ldots, e_n of \mathbb{V} . Then any $\mu \in \operatorname{Hom}(\mathbb{V} \otimes \mathbb{V}, \mathbb{V})$ is determined by n^3 structure constants $c_{ij}^k \in \mathbb{C}$ such that $\mu(e_i \otimes e_j) = \sum_{k=1}^n c_{ij}^k e_k$.

A subset of $\operatorname{Hom}(\mathbb{V} \otimes \mathbb{V}, \mathbb{V})$ is $\operatorname{Zariski-closed}$ if it can be defined by a set of polynomial equations in the variables c_{ij}^k $(1 \leq i, j, k \leq n)$.

Let T be a set of polynomial identities. The set of algebra structures on $\mathbb V$ satisfying polynomial identities from T forms a Zariski-closed subset of the variety $\mathrm{Hom}(\mathbb V\otimes\mathbb V,\mathbb V)$. We denote this subset by $\mathbb L(T)$. The general linear group $GL(\mathbb V)$ acts on $\mathbb L(T)$ by conjugations:

$$(g * \mu)(x \otimes y) = g\mu(g^{-1}x \otimes g^{-1}y)$$

for $x,y\in\mathbb{V},\ \mu\in\mathbb{L}(T)\subset\mathrm{Hom}(\mathbb{V}\otimes\mathbb{V},\mathbb{V})$ and $g\in GL(\mathbb{V})$. Thus, $\mathbb{L}(T)$ is decomposed into $GL(\mathbb{V})$ -orbits that correspond to the isomorphism classes of algebras. Let $O(\mu)$ denote the orbit of $\mu\in\mathbb{L}(T)$ under the action of $GL(\mathbb{V})$ and $\overline{O(\mu)}$ denote the Zariski closure of $O(\mu)$.

Let \mathcal{A} and \mathcal{B} be two n-dimensional algebras satisfying the identities from T, and let $\mu, \lambda \in \mathbb{L}(T)$ represent \mathcal{A} and \mathcal{B} , respectively. We say that \mathcal{A} degenerates to \mathcal{B} and write $\mathcal{A} \to \mathcal{B}$ if $\lambda \in \overline{O(\mu)}$. Note that in this case we have

 $\overline{O(\lambda)} \subset \overline{O(\mu)}$. Hence, the definition of a degeneration does not depend on the choice of μ and λ . If $A \ncong B$, then the assertion $A \to B$ is called a *proper degeneration*. We write $A \not\to B$ if $\lambda \not\in \overline{O(\mu)}$.

Let \mathcal{A} be represented by $\mu \in \mathbb{L}(T)$. Then \mathcal{A} is rigid in $\mathbb{L}(T)$ if $O(\mu)$ is an open subset of $\mathbb{L}(T)$. Recall that a subset of a variety is called irreducible if it cannot be represented as a union of two non-trivial closed subsets. A maximal irreducible closed subset of a variety is called an irreducible component. It is well known that any affine variety can be represented as a finite union of its irreducible components in a unique way. The algebra \mathcal{A} is rigid in $\mathbb{L}(T)$ if and only if $\overline{O(\mu)}$ is an irreducible component of $\mathbb{L}(T)$.

Given the spaces U and W, we write simply U > W instead of dim $U > \dim W$.

2.2. **Method of the description of degenerations of algebras.** In the present work we use the methods applied to Lie algebras in [10, 25, 26, 46]. First of all, if $\mathcal{A} \to \mathcal{B}$ and $\mathcal{A} \ncong \mathcal{B}$, then $\mathfrak{Der}(\mathcal{A}) < \mathfrak{Der}(\mathcal{B})$, where $\mathfrak{Der}(\mathcal{A})$ is the Lie algebra of derivations of \mathcal{A} . We compute the dimensions of algebras of derivations and check the assertion $\mathcal{A} \to \mathcal{B}$ only for such \mathcal{A} and \mathcal{B} that $\mathfrak{Der}(\mathcal{A}) < \mathfrak{Der}(\mathcal{B})$.

To prove degenerations, we construct families of matrices parametrized by t. Namely, let \mathcal{A} and \mathcal{B} be two algebras represented by the structures μ and λ from $\mathbb{L}(T)$ respectively. Let e_1, \ldots, e_n be a basis of \mathbb{V} and c_{ij}^k $(1 \leq i, j, k \leq n)$

be the structure constants of λ in this basis. If there exist $a_i^j(t) \in \mathbb{C}$ $(1 \le i, j \le n, t \in \mathbb{C}^*)$ such that $E_i^t = \sum_{j=1}^n a_i^j(t) e_j$

 $(1 \leq i \leq n)$ form a basis of \mathbb{V} for any $t \in \mathbb{C}^*$, and the structure constants of μ in the basis E_1^t, \ldots, E_n^t are such rational functions $c_{ij}^k(t) \in \mathbb{C}[t]$ that $c_{ij}^k(0) = c_{ij}^k$, then $A \to \mathcal{B}$. In this case E_1^t, \ldots, E_n^t is called a *parametrized basis* for $A \to \mathcal{B}$. To simplify our equations, we will use the notation $A_i = \langle e_i, \ldots, e_n \rangle$, $i = 1, \ldots, n$ and write simply $A_p A_q \subset A_r$ instead of $c_{ij}^k = 0$ $(i \geq p, j \geq q, k \geq r)$.

Since the variety of 4-dimensional nilpotent left symmetric algebras contains infinitely many non-isomorphic algebras, we have to do some additional work. Let $\mathcal{A}(*) := \{\mathcal{A}(\alpha)\}_{\alpha \in I}$ be a series of algebras, and let \mathcal{B} be another algebra. Suppose that for $\alpha \in I$, $\mathcal{A}(\alpha)$ is represented by the structure $\mu(\alpha) \in \mathbb{L}(T)$ and $B \in \mathbb{L}(T)$ is represented by the structure λ . Then we say that $\mathcal{A}(*) \to \mathcal{B}$ if $\lambda \in \overline{\{O(\mu(\alpha))\}_{\alpha \in I}}$, and $\mathcal{A}(*) \to \mathcal{B}$ if $\lambda \notin \overline{\{O(\mu(\alpha))\}_{\alpha \in I}}$.

Let $\mathcal{A}(*)$, \mathcal{B} , $\mu(\alpha)$ ($\alpha \in I$) and λ be as above. To prove $\mathcal{A}(*) \to \mathcal{B}$ it is enough to construct a family of pairs (f(t),g(t)) parametrized by $t \in \mathbb{C}^*$, where $f(t) \in I$ and $g(t) \in GL(\mathbb{V})$. Namely, let e_1,\ldots,e_n be a basis of \mathbb{V} and c_{ij}^k $(1 \le i,j,k \le n)$ be the structure constants of λ in this basis. If we construct $a_i^j:\mathbb{C}^* \to \mathbb{C}$ $(1 \le i,j \le n)$ and

 $f: \mathbb{C}^* \to I$ such that $E_i^t = \sum_{j=1}^n a_i^j(t) e_j$ $(1 \le i \le n)$ form a basis of \mathbb{V} for any $t \in \mathbb{C}^*$, and the structure constants of $\mu_{f(t)}$ in the basis E_1^t, \ldots, E_n^t are such rational functions $c_{ij}^k(t) \in \mathbb{C}[t]$ that $c_{ij}^k(0) = c_{ij}^k$, then $A(*) \to \mathcal{B}$. In this case

 $\mu_{f(t)}$ in the basis E_1^t, \ldots, E_n^t are such rational functions $c_{ij}^v(t) \in \mathbb{C}[t]$ that $c_{ij}^v(0) = c_{ij}^v$, then $\mathcal{A}(*) \to \mathcal{B}$. In this case E_1^t, \ldots, E_n^t and f(t) are called a parametrized basis and a parametrized index for $\mathcal{A}(*) \to \mathcal{B}$, respectively.

We now explain how to prove $\mathcal{A}(*) \not\to \mathcal{B}$. Note that if $\mathfrak{Der} \mathcal{A}(\alpha) > \mathfrak{Der} \mathcal{B}$ for all $\alpha \in I$ then $\mathcal{A}(*) \not\to \mathcal{B}$. One can also use the following Lemma, whose proof is the same as the proof of Lemma 1.5 from [25].

Lemma 5. Let \mathfrak{B} be a Borel subgroup of $GL(\mathbb{V})$ and $\mathfrak{R} \subset \mathbb{L}(T)$ be a \mathfrak{B} -stable closed subset. If $\mathcal{A}(*) \to \mathbb{B}$ and for any $\alpha \in I$ the algebra $\mathcal{A}(\alpha)$ can be represented by a structure $\mu(\alpha) \in \mathbb{R}$, then there is $\lambda \in \mathbb{R}$ representing \mathbb{B} .

2.3. The geometric classification of 4-dimensional nilpotent left symmetric algebras. The main result of the present section is the following theorem.

Theorem B. The variety of 4-dimensional nilpotent left symmetric algebras has dimension 15 and it has three irreducible components defined by infinite families of algebras $\mathbf{L}_{12}^4(\lambda)$, $\mathbf{L}_{21}^4(\lambda)$ and $\mathbf{L}_{23}^4(\lambda,\alpha)$.

Proof. Recall that the description of all irreducible components of 4-dimensional nilpotent Novikov algebras was given in [34]. Using the cited result, we can see that the variety of 4-dimensional Novikov algebras has two irreducible components given by the following families of algebras:

$$\frac{\mathcal{N}_{20}^4(\alpha)}{\mathcal{N}_{22}^4(\lambda)} : e_1e_2 = e_3 \quad e_1e_1 = \alpha e_4 \quad e_1e_3 = e_4 \quad e_2e_2 = e_4 \quad e_2e_3 = e_4 \quad e_3e_2 = -e_4$$

$$\frac{\mathcal{N}_{20}^4(\alpha)}{\mathcal{N}_{22}^4(\lambda)} : e_1e_1 = e_2 \quad e_1e_2 = e_3 \quad e_1e_3 = (2-\lambda)e_4 \quad e_2e_1 = \lambda e_4 \quad e_2e_2 = \lambda e_4 \quad e_3e_1 = \lambda e_4$$

Now we can prove that the variety of 4-dimensional nilpotent left symmetric algebras has three irreducible components. One can easily compute that

$$\mathfrak{Der} \ \mathbf{L}^4_{12}(\lambda) = 2 \quad \mathfrak{Der} \ \mathbf{L}^4_{21}(\lambda) = 2 \quad \mathfrak{Der} \ \mathbf{L}^4_{23}(\lambda,\alpha) = 3$$

The list of all necessary degenerations is given below:

```
E_1^t = \frac{1}{\alpha t - \alpha t^2} e_1 + \frac{1}{\alpha (-1+t)} e_2 + \frac{1+\alpha - 2t - \alpha t}{\alpha^2 (-1+t) t^2} e_3 \quad E_2^t = \frac{1}{\alpha (-1+t)} e_2 + \frac{1}{\alpha^2 t - \alpha^2 t^2} e_3
   \mathbf{L}_{12}^{4}(t)
                                                                                                                                                                                                 E_3^t = \frac{1}{\alpha^2(-1+t)t}e_3 - \frac{1+\alpha}{\alpha^3(-1+t)t^2}e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                         E_3^t = t^3 e_3 E_4^t = -\frac{2t^3}{2t+\lambda} e_4
   \mathbf{L}_{23}^4 \left(t+\lambda, -\frac{2+t}{t^2+t\lambda}\right)
                                                                                                                                                \mathcal{N}^4_{22}(\lambda)
                                                                                                                                                                                                  E_1^t = te_1 E_2^t = t^2 e_2
                                                                                                                                                                                                                                                                                                                                                                                                                                          E_2^t = t^2 e_2 + t^2 e_3 + t e_4
                                                                                                                                                                                                  E_1^t = e_1 + t^2 e_2 + (1 - 2t^2 + t^3)e_3
   \mathbf{L}_{21}^{4}(-t)
                                                                                                                                                                                                                 = te_3 + (1 - t^2 + t^3)e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                          E_4^{\overline{t}} = t^3 e_4
                                                                                                                                                                                                                                                          E_2^t = e_2
                                                                                                                                                                                                                                                                                                                                                                                                                                          E_3^t = t^{-1}e_3
                                                                                                                                                  {f L}_{02}^4
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^t = t^{-1}e_4
  L_{01}^{4}
                                                                                                                                                                                                                                                                                                                                                                                                                                         \frac{E_3 - t}{E_3^t = t^{-1}e_3}
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^t = t^{-3}e_4
                                                                                                                                                 \overline{\mathbf{L}}_{\underline{03}}^{4}
                                                                                                                                                                                                                                                                                                                                                                                                                                                        =t^{-2}e_3
                                                                                                                                                                                                                = t^{-1}e_1 E_2^t = t^{-2}e_2
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                E_4^t = t^{-4}e_4
                                                                                                                                                 \overline{\mathbf{L}}_{\underline{04}}^{4}
 \mathbf{L}_{01}^4
                                                                                                                                                                                                                                                                                                                                                                                                                                         E_3^t = t^{-2}e_3
                                                                                                                                                                                                 E_1^{t} = t^{-1}e_1 	 E_2^{t} = t^{-1}e_2
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                E_4^t = t^{-3}e_4
                                                                                                                                                 L_{05}^{4}
                                                                                                                                                                                                                                                                                                                                                                                                                                         E_2^t = t^4 e_3
                                                                                                                                                                                                E_1^t = t^3 e_1 E_2^t = t^2 e_2
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^{t} = t^5 e_4
  \mathbf{L}_{12}^4(t^{-2})
                                                                                                                                                L_{06}^{4}
                                                                                                                                                                                                E_1^{\tilde{t}} = t^{-1}e_1 E_2^{\tilde{t}} = t^{-1}e_2
                                                                                                                                                                                                                                                                                                                                                                                                                                          E_3^t = t^{-2}e_3
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^t = t^{-3}e_4
                                                                                                                                                 {f L}_{07}^4
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_2^t = -\frac{t^2(-2+2it+t^2)}{(-2+3it+t^2)^2}e_2
                                                                                                                                                                                                                              2t^{3}(2i+2t-it^{2})
                                                                                                                                                                                                 E_1^t = \frac{2t^3(2i+2t-it^2)}{(i+t)^2(2i+t)^3}e_1 + \frac{it^3(-2+2it+t^2)}{(i+t)^2(2i+t)^3}e_2
                                                                                                                                                                                                                                                                                                                                                                                                                               it^4(-2+2it+t^2)^2
   \mathbf{L}_{16}^4(\frac{1-it}{t^2})
                                                                                                                                                                                                                                                                                                                                                                                                                                    (i+t)^4(2i+t)^4
                                                                                                                                                                                                 E_3^t = \frac{t^4(-2+2it+t^2)^2}{(i+t)^4(2i+t)^4}e_3 + \frac{t^4(-2+2it+t^2)^2}{(i+t)^4(2i+t)^4}e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^t = -\frac{2t^4(-2+2it+t^2)^3}{(i+t)^5(2i+t)^6}
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^t = t^{-2}e_4
                                                                                                                                                                                                 E_1^t = t^{-1}e_1 E_2^t = e_2
                                                                                                                                                 L_{09}^{4}
                                                                                                                                                                                                                                                                                                                                                                                                                                          E_3^t = t^{-1}e_3
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^t = \frac{2t^3}{-1+t}e_4
E_4^t = t^{-3}e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                           \overline{E_2^t} = t^2 e_3
   \mathbf{L}_{23}^4(-1, \frac{1+t}{1-t})
                                                                                                                                                \mathbf{L}_{10}^4
                                                                                                                                                                                                 E_1^t = te_1
                                                                                                                                                                                                                                                                      E_2^t = te_2
 \mathbf{L}_{12}^{23}(\lambda)
\mathbf{L}_{14}^{4}(\lambda)
\mathbf{L}_{12}^{4}(t+\lambda)

\begin{array}{ccc}
 & & \mathbf{L}_{11}^{10}(\lambda) \\
 & \rightarrow & & \mathbf{L}_{13}^{4}(\lambda) \\
 & \rightarrow & & \mathbf{L}_{14}^{4}(\lambda)
\end{array}

                                                                                                                                                                                                                                                                                                                                                                                                                                         E_3^t = t^{-2}e_3 
 E_3^t = t^{-2}e_3
                                                                                                                                                                                              E_4^t = t^{-3}e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                           \gamma = \sqrt{1 - 4\lambda}
  E_1^t = -\frac{t(1+\gamma)^2}{2\lambda^2}e_1 + \frac{t(1+\gamma)}{2\lambda^2}e_2 + \frac{t(1+\gamma)}{4\lambda^3(t+\lambda)(1+\gamma)^4 + 4\lambda^2(1+\gamma)^2 \left(t(1-3\lambda) - \lambda^2\right) + 32t^2\lambda^3\right)}{4\lambda^3(t+\lambda)^2(t(1+\gamma) - (1-\gamma)\lambda)}e_3
                                                                                                                                                                                                       4\lambda^3(t+\lambda)^2(t(1+\gamma)-(1-\gamma)\lambda)
E_{2}^{t} = -\frac{t(1-\gamma)}{\lambda^{2}}e_{1} + \frac{2t\left((1-\gamma)(\lambda^{2}-t)+t\lambda(3-\gamma)\right)}{\lambda^{3}(1-\gamma)}e_{2} + \frac{t\left[t^{2}\left(2t(1-\gamma)^{4}-\lambda(1-6\gamma^{2}+8\gamma^{3}-3\gamma^{4}-32\lambda^{2})\right)+(1-\gamma)^{2}\lambda^{2}\left(t(5+2\gamma+\gamma^{2}-12\lambda)-4\lambda^{2}\right)\right]}{2\lambda^{3}(1-\gamma)(t+\lambda)^{2}(t(1-\gamma)-(1+\gamma)\lambda)}e_{3}
E_{3}^{t} = \frac{t^{3}(1+\gamma)^{4}}{4\lambda^{5}}e_{3} - \frac{t^{3}(1+\gamma)^{3}\left(t(t+\lambda)(1+\gamma)^{2}+4\lambda^{2}(1-\lambda-t)\right)}{8\lambda^{6}(t+\lambda)^{2}}e_{4}
E_{4}^{t} = \frac{t^{4}(1+\gamma)^{5}}{8\lambda^{7}}e_{4}
E_{5}^{t} = -\frac{t^{2}(1-\gamma)^{2}}{2\lambda^{2}}e_{1} + \frac{t^{2}(1-\gamma)^{2}}{\lambda^{2}}e_{1} + \frac{t^{2}(1-\gamma)^{4}+4\lambda^{2}(1-\gamma)^{2}\left(t(1-3\lambda)-\lambda^{2}\right)+32t^{2}\lambda^{3}\right)}{4\lambda^{3}(t+\lambda)^{2}\left(t(1-\gamma)-(1+\gamma)\lambda\right)}e_{3}
E_{5}^{t} = -\frac{t(1-\gamma)^{2}}{2\lambda^{2}}e_{1} + \frac{t^{2}(1-\gamma)^{2}}{\lambda^{2}}e_{1} + \frac{t^{2}(1-\gamma)^{4}+4\lambda^{2}(1-\gamma)^{2}\left(t(1-3\lambda)-\lambda^{2}\right)+32t^{2}\lambda^{3}\right)}{4\lambda^{3}(t+\lambda)^{2}\left(t(1-\gamma)-(1+\gamma)\lambda\right)}e_{3}
E_{2}^{t} = -\frac{t(1-\gamma)}{\lambda^{2}}e_{1} + \frac{2t\left((1-\gamma)(\lambda^{2}-t)+t\lambda(3-\gamma)\right)}{\lambda^{3}(1-\gamma)}e_{2} + \frac{t\left[t^{2}\left(2t(1-\gamma)^{4}-\lambda(1-6\gamma^{2}+8\gamma^{3}-3\gamma^{4}-32\lambda^{2})\right)+(1-\gamma)^{2}\lambda^{2}\left(t(5+2\gamma+\gamma^{2}-12\lambda)-4\lambda^{2}\right)\right]}{2\lambda^{3}(1-\gamma)(t+\lambda)^{2}\left(t(1-\gamma)(1-\gamma)(1-\gamma)\lambda\right)}
E_{3}^{t} = \frac{t^{3}(1-\gamma)^{4}}{4\lambda^{5}}e_{3} - \frac{t^{3}(1-\gamma)^{3}\left(t(t+\lambda)(1-\gamma)^{2}+4\lambda^{2}(1-\lambda-t)\right)}{8\lambda^{6}(t+\lambda)^{2}}e_{4} \qquad \qquad E_{4}^{t} = \frac{t^{4}(1-\gamma)^{5}}{8\lambda^{7}}e_{4}
\mathbf{L}_{23}^{t}\left(\frac{1}{t^{3}}, \frac{t^{3}-t^{6}(\alpha-1)}{k^{4}(\alpha-1)}\right) \rightarrow \mathbf{L}_{17}^{t}(\alpha) \quad E_{1}^{t} = te_{1} \qquad E_{2}^{t} = t^{2}e_{2} \qquad \qquad E^{t} - e_{2} \qquad \qquad \mathbf{D}^{t}
                                                                                                                                                                                                                                                                                                                                                                                                                                       E_{3}^{t} = e_{3} \qquad E_{4}^{t} = \frac{2t^{4}}{1+t^{3}(-1+\alpha)}e_{4}
E_{3}^{t} = -\frac{t^{3}(2+t)^{3}}{(1+t)^{3}}e_{3} - \frac{2t^{3}(2+t)^{3}}{(1+t)^{5}}e_{4}
E_{4}^{t} = -\frac{t^{4}(2+t)^{4}}{(1+t)^{5}}e_{4}
                                                                                                                                                                                              E_1^t = -\frac{t(2+t)}{1+t}e_1 - \frac{t(2+t)}{(1+t)^2}e_2
E_2^t = \frac{t^2(2+t)^2}{(1+t)^2}e_2 - \frac{t^2(2+t)^2}{(1+t)^3}e_3 - \frac{t^2(2+t)^2}{(1+t)^3}e_4
  \mathbf{L}_{21}^4(\frac{2}{1+t})
                                                                                                                                               {f L}_{18}^4
                                                                                                                                                                                               E_4^{\hat{t}} = t^{-5}e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                         E_3^t = e_3 - \frac{4t\alpha + 4(1+\alpha) - t^2(1+2\alpha)}{t(6-t+2t\alpha - t^2(1+2\alpha))}e_4
                                                                                                                                                                                              E_1^t = e_1 - \frac{2-t}{t(6-t+2t\alpha-t^2(1+2\alpha))}e_2
                                                                                                                                                \mathbf{L}_{22}^4(\alpha)
   L_{23}^4(t,\frac{2\alpha}{2-t})
                                                                                                                                                                                                                                                                                (2-t)(1+t)
                                                                                                                                                                                                                                                                                                                                                                                                    (2-t)(2+t(2\alpha-1))
                                                                                                                                                                                                E_2^t = e_2 - \frac{(2-t)(1+t)}{t(6-t+2t\alpha-t^2(1+2\alpha))}e_3 - \frac{(2-t)(2+t(2\alpha-1))}{t^2(6-t+2t\alpha-t^2(1+2\alpha))^2}e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                         E_4^t = e_4
   \mathbf{L}_{23}^4(t+\lambda, \frac{3+\lambda}{(\lambda-1)\lambda}) \longrightarrow \mathbf{L}_{24}^4(\lambda)
  \begin{split} & \mathbf{L}_{23}^{t}(t+\lambda,\frac{(\lambda-1)\lambda}{(\lambda-1)\lambda}) &\rightarrow \mathbf{L}_{24}(\lambda) \\ & E_{1}^{t} = e_{1} + \frac{2t^{3}(\lambda^{2}-9)+2t^{2}(9-21\lambda+3\lambda^{2}+\lambda^{3})+t(3+13\lambda-16\lambda^{2}+4\lambda^{3})+2(\lambda-3)}{2t(\lambda-3)(t^{2}(3+\lambda)-t(3-9\lambda-2\lambda^{2})-\lambda(3-6\lambda-\lambda^{2}))} e_{2} \\ & E_{2}^{t} = e_{2} + \frac{2t^{3}(\lambda^{2}-9)+t^{2}(21-29\lambda-10\lambda^{2}+6\lambda^{3})-t(3-18\lambda+3\lambda^{2}+12\lambda^{3}-4\lambda^{4})-2(3+2\lambda-\lambda^{2})}{2t(\lambda-3)(t^{2}(3+\lambda)-t(3-9\lambda-2\lambda^{2})-\lambda(3-6\lambda-\lambda^{2}))} e_{3} \\ & + \left[ \frac{8(\lambda-3)^{2}\lambda(1+\lambda)-2t^{7}(3+\lambda)^{3}(21-13\lambda+2\lambda^{2})+2t^{6}(3+\lambda)^{2}(153-387\lambda+115\lambda^{2}+23\lambda^{3}-8\lambda^{4})+4t(27-27\lambda-93\lambda^{2}+51\lambda^{3}+66\lambda^{4}-48\lambda^{5}+8\lambda^{6})}{4t^{2}\lambda(\lambda-1)(\lambda-3)^{2}(t^{2}(3+\lambda)-t(3-9\lambda-2\lambda^{2})-\lambda(3-6\lambda-\lambda^{2}))^{2}} e_{3} \\ & - \frac{2t^{5}\left[1134-4806\lambda+3789\lambda^{2}+2280\lambda^{3}-894\lambda^{4}-284\lambda^{5}+51\lambda^{6}+10\lambda^{7}\right]+t^{3}\left(27+999\lambda-3591\lambda^{2}+2941\lambda^{3}+368\lambda^{4}-1372\lambda^{5}+432\lambda^{6}+20\lambda^{7}-16\lambda^{8}\right)}{4t^{2}\lambda(\lambda-1)(\lambda-3)^{2}(t^{2}(2+\lambda)-t(2-9\lambda^{2})-\lambda(3-6\lambda-\lambda^{2}))^{2}} e_{3} \\ & - \frac{2t^{5}\left[1134-4806\lambda+3789\lambda^{2}+2280\lambda^{3}-894\lambda^{4}-284\lambda^{5}+51\lambda^{6}+10\lambda^{7}\right]+t^{3}\left(27+999\lambda-3591\lambda^{2}+2941\lambda^{3}+368\lambda^{4}-1372\lambda^{5}+432\lambda^{6}+20\lambda^{7}-16\lambda^{8}\right)}{4t^{2}\lambda(\lambda-1)(\lambda-3)^{2}(t^{2}(2+\lambda)-t(2-9\lambda^{2})-\lambda(3-6\lambda-\lambda^{2}))^{2}} e_{3} \\ & - \frac{2t^{5}\left[1134-4806\lambda+3789\lambda^{2}+2280\lambda^{3}-894\lambda^{4}-284\lambda^{5}+51\lambda^{6}+10\lambda^{7}\right]+t^{3}\left(27+999\lambda-3591\lambda^{2}+2941\lambda^{3}+368\lambda^{4}-1372\lambda^{5}+432\lambda^{6}+20\lambda^{7}-16\lambda^{8}\right)}{4t^{2}\lambda(\lambda-1)(\lambda-3)^{2}(t^{2}(2+\lambda)-t(2-9\lambda^{2})-\lambda(3-6\lambda-\lambda^{2}))^{2}} e_{3} \\ & - \frac{2t^{5}\left[1134-4806\lambda+3789\lambda^{2}+2280\lambda^{3}-894\lambda^{4}-284\lambda^{5}+51\lambda^{6}+10\lambda^{7}\right]+t^{3}\left(27+999\lambda-3591\lambda^{2}+2941\lambda^{3}+368\lambda^{4}-1372\lambda^{5}+432\lambda^{6}+20\lambda^{7}-16\lambda^{8}\right)}{4t^{2}\lambda(\lambda-1)(\lambda-3)^{2}(t^{2}(2+\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t^{2}(2-2\lambda)-t
                                                                                                                                                                                   4t^2\lambda(\lambda-1)(\lambda-3)^2(t^2(3+\lambda)-t(3-9\lambda-2\lambda^2)-\lambda(3-6\lambda-\lambda^2))^2
                +\frac{2t^4 \left(324 - 2700\lambda + 5085\lambda^2 - 1647\lambda^3 - 2218\lambda^4 + 802\lambda^5 + 141\lambda^6 - 39\lambda^7 - 4\lambda^8\right) + 2t^2 \left(54 + 117\lambda - 201\lambda^2 + 63\lambda^3 + 231\lambda^4 + 28\lambda^5 - 228\lambda^6 + 112\lambda^7 - 16\lambda^8\right)}{4t^2\lambda(\lambda - 1)(\lambda - 3)^2(t^2(3 + \lambda) - t(3 - 9\lambda - 2\lambda^2) - \lambda(3 - 6\lambda - \lambda^2))^2}
   +\frac{4t^2\lambda(\lambda-1)(\lambda-3)^2(t^2(3+\lambda)-t(3-9\lambda-2\lambda^2)-\lambda(3-6\lambda-\lambda^2))^2}{4t^2\lambda(\lambda-1)(\lambda-3)^2(t^2(3+\lambda)-t(3-9\lambda-2\lambda^2)-\lambda(3-6\lambda-\lambda^2))^2}\\E_3^t=e_3-\frac{t^3(3+\lambda)^2(4\lambda-13)+t^2\left(144-189\lambda-94\lambda^2+19\lambda^3+8\lambda^4\right)-2t(9-69\lambda+11\lambda^2+37\lambda^3-12\lambda^4\right)-12(3+2\lambda-\lambda^2)}{2t\lambda(t+\lambda)(3-4\lambda+\lambda^2)(3-6\lambda-\lambda^2-t(3+\lambda))}e_4
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 E_4^t = \frac{2(\lambda - 3) + t(3 + \lambda)}{2\lambda(3 - 4\lambda + \lambda^2)} e_4
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REFERENCES

- [1] Abdelwahab H., Calderón A.J., Kaygorodov I., The algebraic and geometric classification of nilpotent binary Lie algebras, International Journal of Algebra and Computation, 29 (2019), 6, 1113–1129.
- [2] Adashev J., Camacho L., Omirov B., Central extensions of null-filiform and naturally graded filiform non-Lie Leibniz algebras, Journal of Algebra, 479 (2017), 461–486.
- [3] Andrada A., Salamon S., Complex product structure on Lie algebras, Forum Mathematicum, 17 (2005), 2, 261–295.
- [4] Bai C., Bijective 1-cocycles and classification of 3-dimensional left-symmetric algebras, Communications in Algebra, 37 (2009), 3, 1016–1057.

- [5] Bakalov B., Kac V., Field algebras, International Mathematics Research Notices. IMRN, (2003), 3, 123–159.
- [6] Balinskii A., Novikov S., Poisson brackets of hydrodynamic type, Frobenius algebras and Lie algebras, Doklady Akademii Nauk SSSR, 32 (1985), 5, 228–231.
- [7] Benes T., Burde D., Classification of orbit closures in the variety of three-dimensional Novikov algebras, Journal of Algebra Appl., 13 (2014), 2, 1350081, 33 pp.
- [8] Bordemann M., Generalized Lax pairs, the modified classical YangBaxter equation, and affine geometry of Lie groups, Communications in Mathematical Physics, 135 (1990), 1, 201–216.
- [9] Burde D., Left-symmetric algebras, or pre-Lie algebras in geometry and physics, Central European Journal of Mathematics, 4 (2006), 3, 323–357.
- [10] Burde D., Steinhoff C., Classification of orbit closures of 4-dimensional complex Lie algebras, Journal of Algebra, 214 (1999), 2, 729–739.
- [11] Cañete E., Khudoyberdiyev A., The classification of 4-dimensional Leibniz algebras, Linear Algebra and its Applications, 439 (2013), 1, 273–288.
- [12] Casas J., Khudoyberdiyev A., Ladra M., Omirov B., On the degenerations of solvable Leibniz algebras, Linear Algebra and its Applications, 439 (2013), 2, 472–487.
- [13] Cayley A., On the theory of analytical forms called trees, Phil. Mag., 13 (1857), 19–30; Collected Math. Papers, University Press, Cambridge, 3 (1890), 242–246.
- [14] Cicalo S., De Graaf W., Schneider C., Six-dimensional nilpotent Lie algebras, Linear Algebra and its Applications, 436 (2012), 1, 163–189.
- [15] Chapoton F., Livernet M., Pre-Lie algebras and the rooted trees operad, International Mathematics Research Notices. IMRN, (2001), 8, 395–408.
- [16] Connes A., Kreimer D., Hopf algebras, renormalization and noncommutative geometry, Communications in Mathematical Physics, 199 (1998), 1, 203–242.
- [17] Darijani I., Usefi H., The classification of 5-dimensional *p*-nilpotent restricted Lie algebras over perfect fields, I., Journal of Algebra, 464 (2016), 97–140.
- [18] De Graaf W., Classification of 6-dimensional nilpotent Lie algebras over fields of characteristic not 2, Journal of Algebra, 309 (2007), 2, 640–653.
- [19] De Graaf W., Classification of nilpotent associative algebras of small dimension, International Journal of Algebra and Computation, 28 (2018), 1, 133–161.
- [20] Etingof P., Schedler T., Soloviev A., Set-theoretical solutions to the quantum YangBaxter equations, Duke Mathematical Journal, 100 (1999), 2, 169–209.
- [21] Gelfand I., Dorfman I., Hamiltonian operators and algebraic structures related to them, Funktsionalnyi Analiz i ego Prilozheniya, 13 (1979), 4, 248–262.
- [22] Gerstenhaber M., The cohomology structure of an associative ring, Annals of Mathematics (2), 78 (1963), 267–288.
- [23] Gorshkov I., Kaygorodov I., Khrypchenko M., The algebraic classification of Tortkara algebras, Communications in Algebra, 2020, DOI: 10.1080/00927872.2020.1742347
- [24] Gorshkov I., Kaygorodov I., Khrypchenko M., The geometric classification of Tortkara algebras, Communications in Algebra, 48 (2020), 1, 204–209.
- [25] Grunewald F., O'Halloran J., Varieties of nilpotent Lie algebras of dimension less than six, Journal of Algebra, 112 (1988), 315–325.
- [26] Grunewald F., O'Halloran J., A Characterization of orbit closure and applications, Journal of Algebra, 116 (1988), 163–175.
- [27] Hegazi A., Abdelwahab H., Classification of five-dimensional nilpotent Jordan algebras, Linear Algebra and its Applications, 494 (2016), 165–218.
- [28] Hegazi A., Abdelwahab H., Calderon Martin A., The classification of n-dimensional non-Lie Malcev algebras with (n-4)-dimensional annihilator, Linear Algebra and its Applications 505 (2016), 32–56.
- [29] Hegazi A., Abdelwahab H., Calderon Martin A., Classification of nilpotent Malcev algebras of small dimensions over arbitrary fields of characteristic not 2, Algebras and Representation Theory, 21 (2018), 1, 19–45.
- [30] Jumaniyozov D., Kaygorodov I., Khudoyberdiyev A., The algebraic and geometric classification of nilpotent noncommutative Jordan algebras, Journal of Algebra and its Applications, to appear, arXiv:1912.02691
- [31] Ismailov N., Kaygorodov I., Mashurov F., The algebraic and geometric classification of nilpotent assosymmetric algebras, Algebras and Representation Theory, 2020, DOI: 10.1007/s10468-019-09935-y
- [32] Ismailov N., Kaygorodov I., Volkov Yu., The geometric classification of Leibniz algebras, International Journal of Mathematics, 29 (2018), 5, 1850035.
- [33] Ismailov N., Kaygorodov I., Volkov Yu., Degenerations of Leibniz and anticommutative algebras, Canadian Mathematical Bulletin, 62 (2019), 3, 539–549.
- [34] Karimjanov I., Kaygorodov I., Khudoyberdiyev A., The algebraic and geometric classification of nilpotent Novikov algebras, Journal of Geometry and Physics, 143 (2019), 11–21.
- [35] Kaygorodov I., Khrypchenko M., Lopes S., The algebraic and geometric classification of nilpotent anticommutative algebras, Journal of Pure and Applied Algebra, 224 (2020), 8, 106337.

- [36] Kaygorodov I., Lopes S., Popov Yu., Degenerations of nilpotent associative commutative algebras, Communications in Algebra, 48 (2020), 4, 1632–1639.
- [37] Kaygorodov I., Páez-Guillán P., Voronin V., The algebraic and geometric classification of nilpotent bicommutative algebras, Algebras and Representation Theory, 2020, DOI: 10.1007/s10468-019-09944-x
- [38] Kaygorodov I., Popov Yu., Pozhidaev A., Volkov Yu., Degenerations of Zinbiel and nilpotent Leibniz algebras, Linear and Multilinear Algebra, 66 (2018), 4, 704–716.
- [39] Kaygorodov I., Popov Yu., Volkov Yu., Degenerations of binary-Lie and nilpotent Malcev algebras, Communications in Algebra, 46 (2018), 11, 4929–4941.
- [40] Kaygorodov I., Volkov Yu., The variety of 2-dimensional algebras over an algebraically closed field, Canadian Journal of Mathematics, 71 (2019), 4, 819–842.
- [41] Kaygorodov I., Volkov Yu., Complete classification of algebras of level two, Moscow Mathematical Journal, 19 (2019), 3, 485–521.
- [42] Kaygorodov I., Volkov Yu., Degenerations of Filippov algebras, Journal of Mathematical Physics, 61 (2020), 2, 021701.
- [43] Khudoyberdiyev A., Rakhimov I., Said Husain Sh.K., On classification of 5-dimensional solvable Leibniz algebras, Linear Algebra and its Applications, 457 (2014), 428–454.
- [44] Koszul J.-L., Domaines bornes homogenes et orbites de groupes de transformations affines, Bulletin de la Societe Mathematique de France, 89 (1961), 515–533.
- [45] Kupershmidt B., Non-abelian phase spaces, Journal of Physics. A. Mathematical and Theoretical, 27 (1994), 8, 2801–2810.
- [46] Seeley C., Degenerations of 6-dimensional nilpotent Lie algebras over C, Communications in Algebra, 18 (1990), 3493–3505.
- [47] Skjelbred T., Sund T., Sur la classification des algebres de Lie nilpotentes, C. R. Acad. Sci. Paris Ser. A-B, 286 (1978), 5, A241–A242.
- [48] Vinberg E., The theory of homogeneous convex cones, Trudy Moskovskogo Matematicheskogo Obshchestva, 12 (1963), 303–358.
- [49] Volkov Yu, Anticommutative Engel algebras of the first five levels, Linear and Multilinear Algebra, 2020, DOI: 10.1080/03081087.2020.1715333
- [50] Zusmanovich P., Central extensions of current algebras, Transactions of the American Mathematical Society, 334 (1992), 1, 143–152.